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Stephen L. Clark

Missouri University of Science and Technology, sclark@mst.edu

Fritz Gesztesy

M. Mitrea

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Boundary Data Maps for Schrödinger Operators on a Compact Interval

S. Clark¹, F. Gesztesy^{2*} and M. Mitrea^{2†}

¹ Department of Mathematics & Statistics, Missouri University of Science and Technology
Rolla, MO 65409, USA

² Department of Mathematics, University of Missouri, Columbia, MO 65211, USA

Dedicated with deep admiration to the memory of Mikhail Sh. Birman (1928-2009)

Abstract. We provide a systematic study of boundary data maps, that is, 2×2 matrix-valued Dirichlet-to-Neumann and more generally, Robin-to-Robin maps, associated with one-dimensional Schrödinger operators on a compact interval $[0, R]$ with separated boundary conditions at 0 and R . Most of our results are formulated in the non-self-adjoint context.

Our principal results include explicit representations of these boundary data maps in terms of the resolvent of the underlying Schrödinger operator and the associated boundary trace maps, Krein-type resolvent formulas relating Schrödinger operators corresponding to different (separated) boundary conditions, and a derivation of the Herglotz property of boundary data maps (up to right multiplication by an appropriate diagonal matrix) in the special self-adjoint case.

Key words: (non-self-adjoint) Schrödinger operators on a compact interval, separated boundary conditions, boundary data maps, Robin-to-Robin maps, linear fractional transformations, Krein-type resolvent formulas

AMS subject classification: 34E05, 34B20, 34L40, 34A55

*Corresponding author. E-mail: gesztesyf@missouri.edu

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1. Introduction

To briefly set the stage for this paper, let $R > 0$, introduce the strip $S_{2\pi} = \{z \in \mathbb{C} \mid 0 \leq \operatorname{Re}(z) < 2\pi\}$, and consider the boundary trace map

$$\gamma_{\theta_0, \theta_R}: \begin{cases} C^1([0, R]) \rightarrow \mathbb{C}^2, \\ u \mapsto \begin{bmatrix} \cos(\theta_0)u(0) + \sin(\theta_0)u'(0) \\ \cos(\theta_R)u(R) - \sin(\theta_R)u'(R) \end{bmatrix}, \end{cases} \quad \theta_0, \theta_R \in S_{2\pi}, \quad (1.1)$$

where “prime” denotes d/dx . In addition, assuming that

$$V \in L^1((0, R); dx) \quad (1.2)$$

(we emphasize that V is not assumed to be real-valued for most of this paper), one can introduce the family of one-dimensional Schrödinger operators H_{θ_0, θ_R} in $L^2((0, R); dx)$ by

$$\begin{aligned} H_{\theta_0, \theta_R} f &= -f'' + Vf, \quad \theta_0, \theta_R \in S_{2\pi}, \\ f \in \operatorname{dom}(H_{\theta_0, \theta_R}) &= \left\{ g \in L^2((0, R); dx) \mid g, g' \in AC([0, R]); \gamma_{\theta_0, \theta_R}(g) = 0; \right. \\ &\quad \left. (-g'' + Vg) \in L^2((0, R); dx) \right\}, \end{aligned} \quad (1.3)$$

where $AC([0, R])$ denotes the set of absolutely continuous functions on $[0, R]$.

Assuming that $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$ (with $\sigma(T)$ denoting the spectrum of T) and $\theta_0, \theta_R \in S_{2\pi}$, we recall that the boundary value problem given by

$$-u'' + Vu = zu, \quad u, u' \in AC([0, R]), \quad (1.4)$$

$$\gamma_{\theta_0, \theta_R}(u) = \begin{bmatrix} c_0 \\ c_R \end{bmatrix} \in \mathbb{C}^2, \quad (1.5)$$

has a unique solution denoted by $u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))$ for each $c_0, c_R \in \mathbb{C}$. To each boundary value problem (1.4), (1.5), we now associate a family of *general boundary data maps*, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) : \mathbb{C}^2 \rightarrow \mathbb{C}^2$, for $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, where

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \begin{bmatrix} c_0 \\ c_R \end{bmatrix} &= \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) (\gamma_{\theta_0, \theta_R}(u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R)))) \\ &= \gamma_{\theta'_0, \theta'_R}(u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))). \end{aligned} \quad (1.6)$$

With $u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))$, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ can be represented as a 2×2 complex matrix, where

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \begin{bmatrix} c_0 \\ c_R \end{bmatrix} &= \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \begin{bmatrix} \cos(\theta_0)u(z, 0) + \sin(\theta_0)u'(z, 0) \\ \cos(\theta_R)u(z, R) - \sin(\theta_R)u'(z, R) \end{bmatrix} \\ &= \begin{bmatrix} \cos(\theta'_0)u(z, 0) + \sin(\theta'_0)u'(z, 0) \\ \cos(\theta'_R)u(z, R) - \sin(\theta'_R)u'(z, R) \end{bmatrix}. \end{aligned} \quad (1.7)$$

The map $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ represents the principal object studied in this paper.

We prove in Section 2. that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ is well-defined for $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$, that is, it is invariant with respect to a change of basis of solutions of (1.4), derive its basic properties (cf. Corollary 2.4), and derive the explicit representation (2.50), (2.51) in terms of a distinguished basis of solutions (2.33).

In Section 3., we relate a special case of $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$, given by the generalized Dirichlet-to-Neumann maps $\Lambda_{\theta_0, \theta_R}(z) = \Lambda_{\theta_0, \theta_R}^{(\theta_0 + \frac{\pi}{2}) \bmod(2\pi), (\theta_R + \frac{\pi}{2}) \bmod(2\pi)}(z)$, to Weyl–Titchmarsh m -functions, derive its asymptotic behavior as z tends to infinity, and most importantly, derive an explicit representation of the boundary data maps $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ in terms of the resolvent of the underlying Schrödinger operator H_{θ_0, θ_R} and the associated boundary trace maps,

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} &= \gamma_{\theta'_0, \theta'_R} [\overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*, \\ \theta_0, \theta_R, \theta'_0, \theta'_R &\in S_{2\pi}, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \end{aligned} \quad (1.8)$$

with $S_{\alpha, \beta}$ denoting the 2×2 diagonal matrix $S_{\alpha, \beta} = \text{diag}(\sin(\alpha), \sin(\beta))$.

Theorem 4.1, the principal result in Section 4., then centers around the following linear fractional transformation relating the boundary data maps $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ and $\Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)$,

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) &= (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} [S_{\delta'_0 - \theta'_0, \delta'_R - \theta'_R} + S_{\theta'_0 - \delta_0, \theta'_R - \delta_R} \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)] \\ &\quad \times [S_{\delta'_0 - \theta_0, \delta'_R - \theta_R} + S_{\theta_0 - \delta_0, \theta_R - \delta_R} \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)]^{-1} S_{\delta'_0 - \delta_0, \delta'_R - \delta_R}, \end{aligned} \quad (1.9)$$

assuming $\theta_0, \theta_R, \theta'_0, \theta'_R, \delta_0, \delta_R, \delta'_0, \delta'_R \in S_{2\pi}$, $\delta'_0 - \delta_0 \neq 0 \bmod(\pi)$, $\delta'_R - \delta_R \neq 0 \bmod(\pi)$, and $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\delta_0, \delta_R}))$. The linear fractional transformation (1.9) then is a major ingredient in our proof that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ is a 2×2 matrix-valued Herglotz function (i.e., analytic on \mathbb{C}_+ , the open complex upper half-plane, with a nonnegative imaginary part) in the special case where H_{θ_0, θ_R} is self-adjoint. (In this case, one necessarily assumes that $\theta_0, \theta_R, \theta'_0, \theta'_R \in [0, 2\pi)$ and that V is real-valued.) In addition, we derive the Herglotz representation of $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ in terms of a 2×2 matrix-valued measure on \mathbb{R} .

The principal result proved in Section 5. then concerns Krein-type resolvent formulas explicitly relating the resolvents of H_{θ_0, θ_R} and $H_{\theta'_0, \theta'_R}$. A typical result to be proved in Theorem 5.3 is of the form

$$\begin{aligned} (H_{\theta'_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} - [\overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} \\ &\quad \times \left[\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right]^{-1} [\gamma_{\theta'_0, \theta'_R} (H_{\theta_0, \theta_R} - zI)^{-1}], \\ \theta_0, \theta_R, \theta'_0, \theta'_R &\in S_{2\pi}, \quad \theta_0 \neq \theta'_0, \quad \theta_R \neq \theta'_R, \quad z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\theta'_0, \theta'_R})). \end{aligned} \quad (1.10)$$

Formula (1.10) demonstrates why $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ is the ideal object for Krein-type resolvent formulas.

Finally, in Section 6., we describe some additional connections between $\Lambda_{\theta_0, \theta_R}(z)$ and the Green's function $G_{\theta_0, \theta_R}(z, x, x')$ of H_{θ_0, θ_R} , and then point out some interesting differences compared to the standard 2×2 Weyl–Titchmarsh matrix associated with H_{θ_0, θ_R} .

For classical as well as recent fundamental literature on Weyl–Titchmarsh operators (i.e., spectral parameter dependent Dirichlet-to-Neumann maps, or more generally, Robin-to-Robin maps, resp., Poincaré–Steklov operators), relevant in the context of boundary value spaces (boundary triples, etc.), we refer, for instance, to [2], [3], [4], [11], [12], [13]–[19], [24]–[28], [33]–[39], [42], [44, Ch. 3], [45], [46, Ch. 13], [49]–[51], [55], [58], [61], [65], [66], [71], [72], [73], [75]–[77], [82] and the references cited therein.

Finally, we briefly summarize some of the notation used in this paper: Let \mathcal{H} be a separable complex Hilbert space, $(\cdot, \cdot)_{\mathcal{H}}$ the scalar product in \mathcal{H} (linear in the second argument), and $I_{\mathcal{H}}$ the identity operator in \mathcal{H} . Next, let T be a linear operator mapping (a subspace of) a Banach space into another, with $\text{dom}(T)$ and $\text{ker}(T)$ denoting the domain and kernel (i.e., null space) of T . The spectrum of a closed linear operator in \mathcal{H} will be denoted by $\sigma(\cdot)$. The Banach space of bounded linear operators on \mathcal{H} is denoted by $\mathcal{B}(\mathcal{H})$, the analogous notation $\mathcal{B}(\mathcal{X}_1, \mathcal{X}_2)$, will be used for bounded operators between two Banach spaces \mathcal{X}_1 and \mathcal{X}_2 . Moreover, $\mathcal{X}_1 \hookrightarrow \mathcal{X}_2$ denotes the continuous embedding of \mathcal{X}_1 into \mathcal{X}_2 .

2. General Boundary Value Problems and Associated Boundary Data Maps

This section is devoted to boundary data maps and their basic properties.

Taking $R > 0$, and fixing $\theta_0, \theta_R \in S_{2\pi}$, with $S_{2\pi}$ the strip

$$S_{2\pi} = \{z \in \mathbb{C} \mid 0 \leq \text{Re}(z) < 2\pi\}, \quad (2.1)$$

we introduce the linear map $\gamma_{\theta_0, \theta_R}$, the trace map associated with the boundary $\{0, R\}$ of $(0, R)$ and the parameters θ_0, θ_R , by

$$\gamma_{\theta_0, \theta_R} : \begin{cases} C^1([0, R]) \rightarrow \mathbb{C}^2, \\ u \mapsto \begin{bmatrix} \cos(\theta_0)u(0) + \sin(\theta_0)u'(0) \\ \cos(\theta_R)u(R) - \sin(\theta_R)u'(R) \end{bmatrix} \end{cases}, \quad \theta_0, \theta_R \in S_{2\pi}, \quad (2.2)$$

where “prime” denotes d/dx . We note, in particular, that the Dirichlet trace γ_D , and the Neumann trace γ_N (in connection with the outward pointing unit normal vector at $\partial(0, R) = \{0, R\}$), are given by

$$\gamma_D = \gamma_{0,0} = -\gamma_{\pi, \pi}, \quad \gamma_N = \gamma_{3\pi/2, 3\pi/2} = -\gamma_{\pi/2, \pi/2}. \quad (2.3)$$

Next, assuming

$$V \in L^1((0, R); dx), \quad (2.4)$$

we introduce the following family of densely defined closed linear operators H_{θ_0, θ_R} acting in

$L^2((0, R); dx)$,

$$\begin{aligned} H_{\theta_0, \theta_R} f &= -f'' + Vf, \quad \theta_0, \theta_R \in S_{2\pi}, \\ f \in \text{dom}(H_{\theta_0, \theta_R}) &= \{g \in L^2((0, R); dx) \mid g, g' \in AC([0, R]); \gamma_{\theta_0, \theta_R}(g) = 0; \\ &\quad (-g'' + Vg) \in L^2((0, R); dx)\}. \end{aligned} \quad (2.5)$$

Here $AC([0, R])$ denotes the set of absolutely continuous functions on $[0, R]$. We emphasize that V is not assumed to be real-valued in the bulk of this paper.

One notices that

$$\gamma_{(\theta_0 + \pi) \bmod(2\pi), (\theta_R + \pi) \bmod(2\pi)} = -\gamma_{\theta_0, \theta_R}, \quad \theta_0, \theta_R \in S_{2\pi}, \quad (2.6)$$

and, on the other hand,

$$H_{(\theta_0 + \pi) \bmod(2\pi), (\theta_R + \pi) \bmod(2\pi)} = H_{\theta_0, \theta_R}, \quad \theta_0, \theta_R \in S_{2\pi}, \quad (2.7)$$

hence it suffices to consider $\theta_0, \theta_R \in S_\pi = \{z \in \mathbb{C} \mid 0 \leq \text{Re}(z) < \pi\}$ rather than $\theta_0, \theta_R \in S_{2\pi}$ in connection with H_{θ_0, θ_R} , but for simplicity of notation we will keep using the strip $S_{2\pi}$ throughout this manuscript.

That H_{θ_0, θ_R} is indeed a closed operator follows, for instance, from [29, Sect. XII.4], especially, by combining Lemma 5 (c) and the first part of the proof of Lemma 26 and noting that $g(0), g'(0)$ (resp., $g(R), g'(R)$) are a complete set of boundary values for the minimal operator H_{\min} associated with the differential expression $-d^2/dx^2 + V(x)$ in $L^2((0, R); dx)$ at $x = 0$ (resp., at $x = R$). Here

$$\begin{aligned} H_{\min} f &= -f'' + Vf, \\ f \in \text{dom}(H_{\min}) &= \{g \in L^2((0, R); dx) \mid g, g' \in AC([0, R]); \\ &\quad g(0) = g'(0) = g(R) = g'(R) = 0; (-g'' + Vg) \in L^2((0, R); dx)\}. \end{aligned} \quad (2.8)$$

Moreover, the adjoint of H_{θ_0, θ_R} is given by

$$\begin{aligned} (H_{\theta_0, \theta_R})^* f &= -f'' + \bar{V}f, \quad \theta_0, \theta_R \in S_{2\pi}, \\ f \in \text{dom}((H_{\theta_0, \theta_R})^*) &= \{g \in L^2((0, R); dx) \mid g, g' \in AC([0, R]); \gamma_{\bar{\theta}_0, \bar{\theta}_R}(g) = 0; \\ &\quad (-g'' + \bar{V}g) \in L^2((0, R); dx)\}. \end{aligned} \quad (2.9)$$

The fact that the spectrum of H_{θ_0, θ_R} , $\sigma(H_{\theta_0, \theta_R})$, is discrete is well-known, but due to its importance in the context of this paper, we now briefly recall its proof following an argument in Marchenko [60]:

Lemma 2.1 (See, [60], Sect. 1.3). *Suppose $V \in L^1((0, R); dx)$, assume $\theta_0, \theta_R \in S_{2\pi}$, and let H_{θ_0, θ_R} be defined as in (2.5). Then $\sigma(H_{\theta_0, \theta_R})$ is an infinite discrete subset of \mathbb{C} (i.e., a set without any finite limit point in \mathbb{C} , but with a limit point at infinity).*

Proof. Fix $z \in \mathbb{C}$. Let $\theta(z, \cdot), \theta'(z, \cdot), \phi(z, \cdot), \phi'(z, \cdot) \in AC([0, R])$, and such that $\theta(z, \cdot)$ and $\phi(z, \cdot)$ are solutions of $-\psi'' + V\psi = z\psi$ uniquely determined by the initial values at $x = 0$,

$$\theta(z, 0) = \phi'(z, 0) = 1, \quad \theta'(z, 0) = \phi(z, 0) = 0. \quad (2.10)$$

Consequently, $\theta(z, \cdot)$ and $\phi(z, \cdot)$ are entire with respect to z . Introducing

$$\psi(z, \cdot) = A\theta(z, \cdot) + B\phi(z, \cdot), \quad A, B \in \mathbb{C}, \quad (2.11)$$

it follows that

$$\gamma_{\theta_0, \theta_R}(\psi) = \begin{bmatrix} \cos(\theta_0)\psi(z, 0) + \sin(\theta_0)\psi'(z, 0) \\ \cos(\theta_R)\psi(z, R) - \sin(\theta_R)\psi'(z, R) \end{bmatrix} = 0 \in \mathbb{C}^2. \quad (2.12)$$

Employing the initial conditions (2.10), one concludes that equation (2.12) is equivalent to

$$\begin{aligned} 0 &= \begin{bmatrix} \cos(\theta_0) & \sin(\theta_0) \\ \cos(\theta_R)\theta(z, R) - \sin(\theta_R)\theta'(z, R) & \cos(\theta_R)\phi(z, R) - \sin(\theta_R)\phi'(z, R) \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} \\ &= \mathcal{U}(z, R, \theta_0, \theta_R) \begin{bmatrix} A \\ B \end{bmatrix}. \end{aligned} \quad (2.13)$$

Consequently, z_0 is an eigenvalue of H_{θ_0, θ_R} if and only if z_0 is a zero of the determinant Δ defined as

$$\Delta(z, R, \theta_0, \theta_R) = \det(\mathcal{U}(z, R, \theta_0, \theta_R)). \quad (2.14)$$

Thus, Δ is an entire function with respect to z , and an explicit computation reveals that

$$\begin{aligned} \Delta(z, R, \theta_0, \theta_R) &= \cos(\theta_0) \cos(\theta_R) \phi(z, R) - \cos(\theta_0) \sin(\theta_R) \phi'(z, R) \\ &\quad - \sin(\theta_0) \cos(\theta_R) \theta(z, R) + \sin(\theta_0) \sin(\theta_R) \theta'(z, R). \end{aligned} \quad (2.15)$$

The standard Volterra integral equations

$$\theta(z, x) = \cos(z^{1/2}x) + \int_0^x dx' \frac{\sin(z^{1/2}(x-x'))}{z^{1/2}} V(x') \theta(z, x'), \quad (2.16)$$

$$\begin{aligned} \phi(z, x) &= \frac{\sin(z^{1/2}x)}{z^{1/2}} + \int_0^x dx' \frac{\sin(z^{1/2}(x-x'))}{z^{1/2}} V(x') \phi(z, x'), \\ &z \in \mathbb{C}, \quad \text{Im}(z^{1/2}) \geq 0, \quad x \in [0, R], \end{aligned} \quad (2.17)$$

then imply that

$$\begin{aligned} \theta(z, x) &\underset{|z| \rightarrow \infty}{=} \cos(z^{1/2}x) + O\left(|z|^{-1/2} e^{\text{Im}(z^{1/2})x}\right), \\ \theta'(z, x) &\underset{|z| \rightarrow \infty}{=} -z^{1/2} \sin(z^{1/2}x) + O\left(e^{\text{Im}(z^{1/2})x}\right), \\ \phi(z, x) &\underset{|z| \rightarrow \infty}{=} \frac{\sin(z^{1/2}x)}{z^{1/2}} + O\left(|z|^{-1} e^{\text{Im}(z^{1/2})x}\right), \\ \phi'(z, x) &\underset{|z| \rightarrow \infty}{=} \cos(z^{1/2}x) + O\left(|z|^{-1/2} e^{\text{Im}(z^{1/2})x}\right). \end{aligned} \quad (2.18)$$

A comparison of (2.15) as $|z| \rightarrow \infty$ and (2.18) demonstrates that Δ does not vanish identically. Thus the set of zeros of Δ , and hence the set of eigenvalues of H_{θ_0, θ_R} , constitutes a discrete set. Again, the asymptotic behavior of Δ near infinity implies that Δ is an entire function of order $1/2$ and hence possesses infinitely many zeros (cf., e.g., [81, p. 252]). \square

In addition (cf. also (3.37), (3.38)), the resolvent of H_{θ_0, θ_R} is clearly a Hilbert–Schmidt operator in $L^2((0, R); dx)$ and the eigenvalues $E_{\theta_0, \theta_R, n}$ of H_{θ_0, θ_R} , in the case of the separated boundary conditions at hand, are of the form $E_{\theta_0, \theta_R, n} = [(n\pi/R) + (a_n/n)]^2$ with $a_n \in \ell^\infty(\mathbb{N})$ as $n \rightarrow \infty$, as shown in [60, Lemma 1.3.3].

Having described the operator H_{θ_0, θ_R} in some detail, still assuming (2.4), we now briefly recall the corresponding closed, sectorial, and densely defined sesquilinear form, denoted by $Q_{H_{\theta_0, \theta_R}}$, associated with H_{θ_0, θ_R} (cf. [47, p. 312, 321, 327–328]):

$$Q_{H_{\theta_0, \theta_R}}(f, g) = \int_0^R dx [f'(x)\overline{g'(x)} + V(x)\overline{f(x)}g(x)] - \cot(\theta_0)\overline{f(0)}g(0) - \cot(\theta_R)\overline{f(R)}g(R), \quad (2.19)$$

$$\begin{aligned} f, g \in \text{dom}(Q_{H_{\theta_0, \theta_R}}) &= H^1((0, R)) \\ &= \{h \in L^2((0, R); dx) \mid h \in AC([0, R]); h' \in L^2((0, R); dx)\}, \\ &\quad \theta_0, \theta_R \in S_{2\pi} \setminus \{0, \pi\}, \end{aligned}$$

$$Q_{H_{0, \theta_R}}(f, g) = \int_0^R dx [f'(x)\overline{g'(x)} + V(x)\overline{f(x)}g(x)] - \cot(\theta_R)\overline{f(R)}g(R), \quad (2.20)$$

$$\begin{aligned} f, g \in \text{dom}(Q_{H_{0, \theta_R}}) & \\ &= \{h \in L^2((0, R); dx) \mid h \in AC([0, R]); h(0) = 0; h' \in L^2((0, R); dx)\}, \\ &\quad \theta_R \in S_{2\pi} \setminus \{0, \pi\}, \end{aligned}$$

$$Q_{H_{\theta_0, 0}}(f, g) = \int_0^R dx [f'(x)\overline{g'(x)} + V(x)\overline{f(x)}g(x)] - \cot(\theta_0)\overline{f(0)}g(0), \quad (2.21)$$

$$\begin{aligned} f, g \in \text{dom}(Q_{H_{\theta_0, 0}}) & \\ &= \{h \in L^2((0, R); dx) \mid h \in AC([0, R]); h(R) = 0; h' \in L^2((0, R); dx)\}, \\ &\quad \theta_0 \in S_{2\pi} \setminus \{0, \pi\}, \end{aligned}$$

$$Q_{H_{0, 0}}(f, g) = \int_0^R dx [f'(x)\overline{g'(x)} + V(x)\overline{f(x)}g(x)], \quad (2.22)$$

$$\begin{aligned} f, g \in \text{dom}(Q_{H_{0, 0}}) &= \text{dom}(|H_{0, 0}|^{1/2}) = H_0^1((0, R)) \\ &= \{h \in L^2((0, R); dx) \mid h \in AC([0, R]); h(0) = 0, h(R) = 0; \\ &\quad h' \in L^2((0, R); dx)\}. \end{aligned}$$

Equations (2.19)–(2.22) follow from the fact that for any $\varepsilon > 0$, there exists $\eta(\varepsilon) > 0$ such that

for all $h \in H^1((0, R))$,

$$|h(x_0)| \leq \varepsilon \|h'\|_{L^2((0,R);dx)} + \eta(\varepsilon) \|h\|_{L^2((0,R);dx)}, \quad x_0 \in [0, R], \quad (2.23)$$

$$\| |V|^{1/2} h \|_{L^2((0,R);dx)} \leq \varepsilon \|h'\|_{L^2((0,R);dx)} + \eta(\varepsilon) \|h\|_{L^2((0,R);dx)} \quad (2.24)$$

(cf. [47, p. 193, 345–346]).

Next, we recall the following elementary, yet fundamental, fact:

Lemma 2.2. *Suppose that $V \in L^1((0, R); dx)$, fix $\theta_0, \theta_R \in S_{2\pi}$, and assume that $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$. Then the boundary value problem given by*

$$-u'' + Vu = zu, \quad u, u' \in AC([0, R]), \quad (2.25)$$

$$\gamma_{\theta_0, \theta_R}(u) = \begin{bmatrix} c_0 \\ c_R \end{bmatrix} \in \mathbb{C}^2, \quad (2.26)$$

has a unique solution $u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))$ for each $c_0, c_R \in \mathbb{C}$.

Proof. This is well-known, but for the sake of completeness, we briefly recall the argument: Let $\psi_j(z, \cdot)$, $j = 1, 2$, be a basis for the solutions of (2.25) and let $\psi(z, \cdot) = A\psi_1(z, \cdot) + B\psi_2(z, \cdot)$, $A, B \in \mathbb{C}$, be the general solution of (2.25). Then

$$\begin{aligned} \gamma_{\theta_0, \theta_R}(\psi(z, \cdot)) &= \begin{bmatrix} \gamma_{\theta_0, \theta_R}(\psi_1(z, \cdot)) & \gamma_{\theta_0, \theta_R}(\psi_2(z, \cdot)) \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} \\ &= \begin{bmatrix} M_{1,1}(z) & M_{1,2}(z) \\ M_{2,1}(z) & M_{2,2}(z) \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix}, \end{aligned} \quad (2.27)$$

where

$$\begin{aligned} M_{1,1}(z) &= \cos(\theta_0)\psi_1(z, 0) + \sin(\theta_0)\psi_1'(z, 0), \\ M_{1,2}(z) &= \cos(\theta_0)\psi_2(z, 0) + \sin(\theta_0)\psi_2'(z, 0), \\ M_{2,1}(z) &= \cos(\theta_R)\psi_1(z, R) - \sin(\theta_R)\psi_1'(z, R), \\ M_{2,2}(z) &= \cos(\theta_R)\psi_2(z, R) - \sin(\theta_R)\psi_2'(z, R). \end{aligned} \quad (2.28)$$

Thus, prescribing $c_0, c_R \in \mathbb{C}$, the equation

$$\gamma_{\theta_0, \theta_R}(\psi(z, \cdot)) = \begin{bmatrix} c_0 \\ c_R \end{bmatrix} \quad (2.29)$$

is uniquely solvable in terms of some $A, B \in \mathbb{C}$ if and only if

$$\det \left(\begin{bmatrix} \gamma_{\theta_0, \theta_R}(\psi_1(z, \cdot)) & \gamma_{\theta_0, \theta_R}(\psi_2(z, \cdot)) \end{bmatrix} \right) = \det \left(\begin{bmatrix} M_{1,1}(z) & M_{1,2}(z) \\ M_{2,1}(z) & M_{2,2}(z) \end{bmatrix} \right) \neq 0. \quad (2.30)$$

On the other hand, this determinant equals zero for some $z_0 \in \mathbb{C}$, if and only if there is a nonzero vector $\begin{bmatrix} A_0 & B_0 \end{bmatrix}^\top \in \mathbb{C}^2$ such that

$$\begin{bmatrix} \gamma_{\theta_0, \theta_R}(\psi_1(z_0, \cdot)) & \gamma_{\theta_0, \theta_R}(\psi_2(z_0, \cdot)) \end{bmatrix} \begin{bmatrix} A_0 \\ B_0 \end{bmatrix} = 0 \quad (2.31)$$

which is equivalent to the existence of a nonzero solution $\psi_0(z_0, \cdot) = A_0\psi_1(z_0, \cdot) + B_0\psi_2(z_0, \cdot)$ of the corresponding boundary value problem given by (2.25) and (2.26) with $z = z_0$ and homogeneous boundary conditions (i.e., with $c_0 = c_R = 0$). Equivalently, $\psi_0(z_0, \cdot)$ satisfies

$$H_{\theta_0, \theta_R}\psi_0(z_0, \cdot) = z_0\psi_0(z_0, \cdot), \quad \psi_0(z_0, \cdot) \in \text{dom}(H_{\theta_0, \theta_R}), \quad (2.32)$$

which in turn is equivalent to $z_0 \in \sigma(H_{\theta_0, \theta_R})$. \square

Assuming $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$, a basis for the solutions of (2.25) is given by

$$\begin{aligned} u_{-, \theta_0}(z, \cdot) &= u(z, \cdot; (\theta_0, 0), (0, 1)), \\ u_{+, \theta_R}(z, \cdot) &= u(z, \cdot; (0, 1), (\theta_R, 0)). \end{aligned} \quad (2.33)$$

Explicitly, one then has

$$u_{-, \theta_0}(z, R) = 1, \quad \cos(\theta_0)u_{-, \theta_0}(z, 0) + \sin(\theta_0)u'_{-, \theta_0}(z, 0) = 0, \quad (2.34)$$

$$u_{+, \theta_R}(z, 0) = 1, \quad \cos(\theta_R)u_{+, \theta_R}(z, R) - \sin(\theta_R)u'_{+, \theta_R}(z, R) = 0. \quad (2.35)$$

Recalling the Wronskian of two functions f and g ,

$$W(f, g)(x) = f(x)g'(x) - f'(x)g(x), \quad f, g \in C^1([0, R]), \quad (2.36)$$

one then computes

$$\begin{aligned} &W(u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot)) \\ &= u_{+, \theta_R}(z, x)u'_{-, \theta_0}(z, x) - u'_{+, \theta_R}(z, x)u_{-, \theta_0}(z, x) \neq 0, \quad x \in [0, R], \\ &= u'_{-, \theta_0}(z, 0) - u'_{+, \theta_R}(z, 0), u_{-, \theta_0}(z, 0) \end{aligned} \quad (2.37)$$

$$= u_{+, \theta_R}(z, R)u'_{-, \theta_0}(z, R) - u'_{+, \theta_R}(z, R). \quad (2.38)$$

To each boundary value problem (2.25), (2.26), we now associate a family of *general boundary data maps*, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) : \mathbb{C}^2 \rightarrow \mathbb{C}^2$, for $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, where

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \begin{bmatrix} c_0 \\ c_R \end{bmatrix} &= \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) (\gamma_{\theta_0, \theta_R}(u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R)))) \\ &= \gamma_{\theta'_0, \theta'_R}(u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))). \end{aligned} \quad (2.39)$$

With $u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))$, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ can be represented as a 2×2 complex matrix, where

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \begin{bmatrix} c_0 \\ c_R \end{bmatrix} &= \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \begin{bmatrix} \cos(\theta_0)u(z, 0) + \sin(\theta_0)u'(z, 0) \\ \cos(\theta_R)u(z, R) - \sin(\theta_R)u'(z, R) \end{bmatrix} \\ &= \begin{bmatrix} \cos(\theta'_0)u(z, 0) + \sin(\theta'_0)u'(z, 0) \\ \cos(\theta'_R)u(z, R) - \sin(\theta'_R)u'(z, R) \end{bmatrix}. \end{aligned} \quad (2.40)$$

The following result shows that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ is well-defined for $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$, that is, it is invariant with respect to a change of basis of solutions of (2.25).

Theorem 2.3. *Let $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$ and $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$. In addition, denote by $\psi_j(z, \cdot)$, $j = 1, 2$, a basis for the solutions of (2.25). Then,*

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) & \quad (2.41) \\ &= \begin{bmatrix} \cos(\theta'_0)\psi_1(z, 0) + \sin(\theta'_0)\psi'_1(z, 0) & \cos(\theta'_0)\psi_2(z, 0) + \sin(\theta'_0)\psi'_2(z, 0) \\ \cos(\theta'_R)\psi_1(z, R) - \sin(\theta'_R)\psi'_1(z, R) & \cos(\theta'_R)\psi_2(z, R) - \sin(\theta'_R)\psi'_2(z, R) \end{bmatrix} \\ & \quad \times \begin{bmatrix} \cos(\theta_0)\psi_1(z, 0) + \sin(\theta_0)\psi'_1(z, 0) & \cos(\theta_0)\psi_2(z, 0) + \sin(\theta_0)\psi'_2(z, 0) \\ \cos(\theta_R)\psi_1(z, R) - \sin(\theta_R)\psi'_1(z, R) & \cos(\theta_R)\psi_2(z, R) - \sin(\theta_R)\psi'_2(z, R) \end{bmatrix}^{-1}. \end{aligned}$$

Moreover, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ is invariant with respect to a change of basis for the solutions of (2.25).

Proof. Letting $\psi(z, \cdot) = A\psi_1(z, \cdot) + B\psi_2(z, \cdot)$, $A, B \in \mathbb{C}$, be an arbitrary solution of (2.25), one observes, by (2.27) and (2.40), that the equation $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\gamma_{\theta_0, \theta_R}(\psi)) = \gamma_{\theta'_0, \theta'_R}(\psi)$ becomes

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) & \quad (2.42) \\ & \times \begin{bmatrix} \cos(\theta_0)\psi_1(z, 0) + \sin(\theta_0)\psi'_1(z, 0) & \cos(\theta_0)\psi_2(z, 0) + \sin(\theta_0)\psi'_2(z, 0) \\ \cos(\theta_R)\psi_1(z, R) - \sin(\theta_R)\psi'_1(z, R) & \cos(\theta_R)\psi_2(z, R) - \sin(\theta_R)\psi'_2(z, R) \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} \\ &= \begin{bmatrix} \cos(\theta'_0)\psi_1(z, 0) + \sin(\theta'_0)\psi'_1(z, 0) & \cos(\theta'_0)\psi_2(z, 0) + \sin(\theta'_0)\psi'_2(z, 0) \\ \cos(\theta'_R)\psi_1(z, R) - \sin(\theta'_R)\psi'_1(z, R) & \cos(\theta'_R)\psi_2(z, R) - \sin(\theta'_R)\psi'_2(z, R) \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} \end{aligned}$$

for every $\begin{bmatrix} A & B \end{bmatrix}^\top \in \mathbb{C}^2$. By the invertibility of $\begin{bmatrix} \gamma_{\theta_0, \theta_R}(\psi_1) & \gamma_{\theta_0, \theta_R}(\psi_2) \end{bmatrix}$ noted in (2.30), equation (2.41) then follows.

Let $\phi_j(z, \cdot)$, $j = 1, 2$, denote a second basis for the solutions of (2.25). Then, there is a nonsingular matrix $K \in \mathbb{C}^{2 \times 2}$ such that $\begin{bmatrix} \psi_1 & \psi_2 \end{bmatrix} = \begin{bmatrix} \phi_1 & \phi_2 \end{bmatrix} K$. Next, we introduce for each pair, $\theta_0, \theta_R \in S_{2\pi}$, the following matrices

$$C_{\theta_0, \theta_R} = \begin{bmatrix} \cos(\theta_0) & 0 \\ 0 & \cos(\theta_R) \end{bmatrix}, \quad S_{\theta_0, \theta_R} = \begin{bmatrix} \sin(\theta_0) & 0 \\ 0 & \sin(\theta_R) \end{bmatrix}. \quad (2.43)$$

Introducing θ_j to denote θ_0 , and θ_k to denote θ_R , respectively; or, using θ_j to denote θ'_0 and θ_k to denote θ'_R , one computes,

$$\begin{aligned} & \begin{bmatrix} \gamma_{\theta_j, \theta_k}(\psi_1(z, \cdot)) & \gamma_{\theta_j, \theta_k}(\psi_2(z, \cdot)) \end{bmatrix} \\ &= C_{\theta_j, \theta_k} \begin{bmatrix} \psi_1(z, 0) & \psi_2(z, 0) \\ \psi_1(z, R) & \psi_2(z, R) \end{bmatrix} + S_{\theta_j, \theta_k} \begin{bmatrix} \psi'_1(z, 0) & \psi'_2(z, 0) \\ -\psi'_1(z, R) & -\psi'_2(z, R) \end{bmatrix} \\ &= \left(C_{\theta_j, \theta_k} \begin{bmatrix} \phi_1(z, 0) & \phi_2(z, 0) \\ \phi_1(z, R) & \phi_2(z, R) \end{bmatrix} + S_{\theta_j, \theta_k} \begin{bmatrix} \phi'_1(z, 0) & \phi'_2(z, 0) \\ -\phi'_1(z, R) & -\phi'_2(z, R) \end{bmatrix} \right) K \\ &= \begin{bmatrix} \gamma_{\theta_j, \theta_k}(\phi_1(z, \cdot)) & \gamma_{\theta_j, \theta_k}(\phi_2(z, \cdot)) \end{bmatrix} K. \end{aligned} \quad (2.44)$$

As defined in (2.41), $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} = \begin{bmatrix} \gamma_{\theta'_0, \theta'_R}(\psi_1) & \gamma_{\theta'_0, \theta'_R}(\psi_2) \end{bmatrix} \begin{bmatrix} \gamma_{\theta_0, \theta_R}(\psi_1) & \gamma_{\theta_0, \theta_R}(\psi_2) \end{bmatrix}^{-1}$, and by (2.44),

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} &= \begin{bmatrix} \gamma_{\theta'_0, \theta'_R}(\psi_1) & \gamma_{\theta'_0, \theta'_R}(\psi_2) \end{bmatrix} \begin{bmatrix} \gamma_{\theta_0, \theta_R}(\psi_1) & \gamma_{\theta_0, \theta_R}(\psi_2) \end{bmatrix}^{-1} \\ &= \begin{bmatrix} \gamma_{\theta'_0, \theta'_R}(\phi_1) & \gamma_{\theta'_0, \theta'_R}(\phi_2) \end{bmatrix} \begin{bmatrix} \gamma_{\theta_0, \theta_R}(\phi_1) & \gamma_{\theta_0, \theta_R}(\phi_2) \end{bmatrix}^{-1}, \end{aligned} \quad (2.45)$$

completing the proof. \square

Theorem 2.3 then readily implies the following result:

Corollary 2.4. *Let $\theta_0, \theta_R, \theta'_0, \theta'_R, \theta''_0, \theta''_R \in S_{2\pi}$. Then, with I_2 denoting the identity matrix in \mathbb{C}^2 ,*

$$\Lambda_{\theta_0, \theta_R}^{\theta_0, \theta_R}(z) = I_2, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad (2.46)$$

$$\Lambda_{\theta'_0, \theta'_R}^{\theta'_0, \theta'_R}(z) \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) = \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z), \quad z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\theta'_0, \theta'_R})), \quad (2.47)$$

$$\Lambda_{\theta'_0, \theta'_R}^{\theta_0, \theta_R}(z) = \left[\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right]^{-1}, \quad z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\theta'_0, \theta'_R})). \quad (2.48)$$

Remark 2.5. *By Theorem 2.3, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ is invariant with respect to a change of basis for the solutions of (2.25). However, the representation of $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ with respect to a specific basis can be simplified considerably with an appropriate choice of basis. For example, by choosing the basis given in (2.33), and by letting $\psi_1(z, \cdot) = u_{+, \theta_R}(z, \cdot) = u(z, \cdot; (0, 1), (\theta_R, 0))$, and $\psi_2(z, \cdot) = u_{-, \theta_0}(z, \cdot) = u(z, \cdot; (\theta_0, 0), (0, 1))$, (2.45) implies that, entrywise,*

$$\begin{aligned} & \left[\gamma_{\theta_0, \theta_R}(u_{+, \theta_R}(z, \cdot)) \quad \gamma_{\theta_0, \theta_R}(u_{-, \theta_0}(z, \cdot)) \right]_{1,1} \\ & \quad = \cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0), \\ & \left[\gamma_{\theta_0, \theta_R}(u_{+, \theta_R}(z, \cdot)) \quad \gamma_{\theta_0, \theta_R}(u_{-, \theta_0}(z, \cdot)) \right]_{1,2} \\ & \quad = \cos(\theta_0)u_{-, \theta_0}(z, 0) + \sin(\theta_0)u'_{-, \theta_0}(z, 0), \\ & \left[\gamma_{\theta_0, \theta_R}(u_{+, \theta_R}(z, \cdot)) \quad \gamma_{\theta_0, \theta_R}(u_{-, \theta_0}(z, \cdot)) \right]_{2,1} \\ & \quad = \cos(\theta_R)u_{+, \theta_R}(z, R) - \sin(\theta_R)u'_{+, \theta_R}(z, R), \\ & \left[\gamma_{\theta_0, \theta_R}(u_{+, \theta_R}(z, \cdot)) \quad \gamma_{\theta_0, \theta_R}(u_{-, \theta_0}(z, \cdot)) \right]_{2,1} \\ & \quad = \cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R). \end{aligned} \quad (2.49)$$

Hence, for this basis,

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) = \left[\left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{j,k} \right]_{1 \leq j, k \leq 2}, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad (2.50)$$

$$\begin{aligned} \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{1,1} &= \frac{\cos(\theta'_0) + \sin(\theta'_0)u'_{+, \theta_R}(z, 0)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)}, \\ \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{1,2} &= \frac{\cos(\theta'_0)u_{-, \theta_0}(z, 0) + \sin(\theta'_0)u'_{-, \theta_0}(z, 0)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)}, \\ \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{2,1} &= \frac{\cos(\theta'_R)u_{+, \theta_R}(z, R) - \sin(\theta'_R)u'_{+, \theta_R}(z, R)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)}, \\ \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{2,2} &= \frac{\cos(\theta'_R) - \sin(\theta'_R)u'_{-, \theta_0}(z, R)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)}. \end{aligned} \quad (2.51)$$

In particular, by (2.34) and (2.35),

$$\left(\Lambda_{\theta_0, \theta'_R}^{\theta_0, \theta'_R}(z)\right)_{1,2} = 0, \quad \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta_R}(z)\right)_{2,1} = 0. \quad (2.52)$$

Remark 2.6. We note that $\Lambda_{0,0}^{\frac{\pi}{2}, \frac{\pi}{2}}(z)$ represents the Dirichlet-to-Neumann map, $\Lambda_{D,N}(z)$, for the boundary value problem (2.25), (2.26); that is, when $\theta_0 = \theta_R = 0$, $\theta'_0 = \theta'_R = \pi/2$, then (2.40) becomes

$$\Lambda_{D,N}(z) \begin{bmatrix} u(z, 0) \\ u(z, R) \end{bmatrix} = \Lambda_{0,0}^{\frac{\pi}{2}, \frac{\pi}{2}}(z) \begin{bmatrix} u(z, 0) \\ u(z, R) \end{bmatrix} = \begin{bmatrix} u'(z, 0) \\ -u'(z, R) \end{bmatrix}, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad (2.53)$$

with $u(z, \cdot) = u(z, \cdot; (0, c_0), (0, c_R))$, $u(z, 0) = c_0$, $u(z, R) = c_R$. The Dirichlet-to-Neumann map in the case $V = 0$ has recently been considered in [72, Example 5.1]. The Neumann-to-Dirichlet map $\Lambda_{N,D}(z) = \Lambda_{\frac{\pi}{2}, \pi/2}^{\pi, \pi}(z) = -[\Lambda_{D,N}(z)]^{-1}$ (cf. (4.70)) in the case $V = 0$ has earlier been computed in [27, Example 4.1]. We also refer to [12], [19], [26], [30] in the intimately related context of Q and M -functions.

It would be interesting to establish precise connections between $\Lambda_{\theta_0, \theta_R}^{\theta_0, \theta_R}(z)$ and the dynamical response operator discussed, for instance, in [7], [8], [9], [10] in connection with the problem of regularity and controllability of the wave equation on a compact interval.

We conclude this section with an elementary result needed in the proof of Lemma 3.4 and Theorem 4.6:

Lemma 2.7. Let $V \in L^1((0, R); dx)$, fix $\theta_0, \theta_R \in S_{2\pi}$, $c_0, c_R \in \mathbb{C}$, and assume that $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$. Then the unique solution $u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R))$ of (2.25), (2.26) can be expressed in terms of the fundamental system $\theta(z, \cdot), \phi(z, \cdot)$, as introduced in the proof of Lemma 2.1, and satisfying the initial conditions (2.10), as follows:

$$u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R)) = \alpha(z)\theta(z, \cdot) + \beta(z)\phi(z, \cdot), \quad (2.54)$$

$$\begin{bmatrix} \alpha(z) \\ \beta(z) \end{bmatrix} = \frac{1}{\Delta(z, R, \theta_0, \theta_R)} \begin{bmatrix} [\cos(\theta_R)\phi(z, R) - \sin(\theta_R)\phi'(z, R)]c_0 - \sin(\theta_0)c_R \\ -[\cos(\theta_R)\theta(z, R) - \sin(\theta_R)\theta'(z, R)]c_0 + \cos(\theta_0)c_R \end{bmatrix},$$

with Δ given by (2.15). In particular, one has

$$u_{+, \theta_R}(z, \cdot) = u(z, \cdot; ((0, 1), (\theta_R, 0))) = \alpha(z)\theta(z, \cdot) + \beta(z)\phi(z, \cdot),$$

$$\begin{bmatrix} \alpha(z) \\ \beta(z) \end{bmatrix} = \frac{1}{\Delta(z, R, 0, \theta_R)} \begin{bmatrix} \cos(\theta_R)\phi(z, R) - \sin(\theta_R)\phi'(z, R) \\ -\cos(\theta_R)\theta(z, R) + \sin(\theta_R)\theta'(z, R) \end{bmatrix}, \quad (2.55)$$

and

$$u_{-, \theta_0}(z, \cdot) = u(z, \cdot; (\theta_0, 0), (0, 1)) \quad (2.56)$$

$$= \frac{1}{\Delta(z, R, \theta_0, 0)} (-\sin(\theta_0)\theta(z, \cdot) + \cos(\theta_0)\phi(z, \cdot)).$$

Proof. With $u(z, \cdot) = u(z, \cdot; (\theta_0, c_0), (\theta_R, c_R)) = \alpha\theta(z, \cdot) + \beta\phi(z, \cdot)$ one observes from (2.26) that

$$\begin{aligned}
\begin{bmatrix} c_0 \\ c_R \end{bmatrix} &= \gamma_{\theta_0, \theta_R}(u(z, \cdot)) = \gamma_{\theta_0, \theta_R}(\alpha\theta(z, \cdot) + \beta\phi(z, \cdot)) \\
&= [\gamma_{\theta_0, \theta_R}(\theta(z, \cdot)) \quad \gamma_{\theta_0, \theta_R}(\phi(z, \cdot))] \begin{bmatrix} \alpha \\ \beta \end{bmatrix} \\
&= \begin{bmatrix} \cos(\theta_0) & \sin(\theta_0) \\ \cos(\theta_R)\theta(z, R) - \sin(\theta_R)\theta'(z, R) & \cos(\theta_R)\phi(z, R) - \sin(\theta_R)\phi'(z, R) \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix} \\
&= \mathcal{U}(z, R, \theta_0, \theta_R) \begin{bmatrix} \alpha \\ \beta \end{bmatrix}, \tag{2.57}
\end{aligned}$$

with \mathcal{U} introduced in (2.13). Solving (2.57) for α, β (and observing that $\det(\mathcal{U}) = \Delta$ according to (2.14)) then yields (2.54) and hence the special cases (2.55) and (2.56). \square

3. Resolvent Formulas

In the principal result of this section, we will derive a formula for $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ in terms of the resolvent of H_{θ_0, θ_R} and the boundary traces $\gamma_{\theta'_0, \theta'_R}$. But first we focus on the special boundary data map given by

$$\Lambda_{\theta_0, \theta_R}(z) = \Lambda_{\theta_0, \theta_R}^{(\theta_0 + \frac{\pi}{2}) \bmod(2\pi), (\theta_R + \frac{\pi}{2}) \bmod(2\pi)}(z), \quad \theta_0, \theta_R \in S_{2\pi}, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \tag{3.1}$$

a generalization of the Dirichlet-to-Neumann map $\Lambda_{D,N}(z) = \Lambda_{0,0}(z)$. Using the basis for solutions of (2.25) given in (2.33) by $u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot)$, and with $\theta'_0 = (\theta_0 + \pi/2) \bmod(2\pi), \theta'_R = (\theta_R + \pi/2) \bmod(2\pi)$, equation (2.50) becomes

$$\Lambda_{\theta_0, \theta_R}(z) = [(\Lambda_{\theta_0, \theta_R}(z))_{j,k}]_{1 \leq j, k \leq 2}, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \tag{3.2}$$

$$\begin{aligned}
(\Lambda_{\theta_0, \theta_R}(z))_{1,1} &= \frac{-\sin(\theta_0) + \cos(\theta_0)u'_{+, \theta_R}(z, 0)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)}, \\
(\Lambda_{\theta_0, \theta_R}(z))_{1,2} &= \frac{-\sin(\theta_0)u_{-, \theta_0}(z, 0) + \cos(\theta_0)u'_{-, \theta_0}(z, 0)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)}, \\
(\Lambda_{\theta_0, \theta_R}(z))_{2,1} &= \frac{-\sin(\theta_R)u_{+, \theta_R}(z, R) - \cos(\theta_R)u'_{+, \theta_R}(z, R)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)}, \\
(\Lambda_{\theta_0, \theta_R}(z))_{2,2} &= \frac{-\sin(\theta_R) - \cos(\theta_R)u'_{-, \theta_0}(z, R)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)}. \tag{3.3}
\end{aligned}$$

Next, consider

$$\alpha(\xi) = [\cos(\xi) \quad \sin(\xi)], \quad \xi \in S_{2\pi}, \tag{3.4}$$

and let $\Psi(z, \cdot; x_0, \alpha(\xi))$ denote the fundamental matrix of solutions of (2.25) given by

$$\begin{aligned}\Psi(z, \cdot; x_0, \alpha(\xi)) &= \begin{bmatrix} \Theta(z, \cdot; x_0, \alpha(\xi)) & \Phi(z, \cdot; x_0, \alpha(\xi)) \\ \vartheta(z, \cdot; x_0, \alpha(\xi)) & \varphi(z, \cdot; x_0, \alpha(\xi)) \end{bmatrix} \\ &= \begin{bmatrix} \vartheta(z, \cdot; x_0, \alpha(\xi)) & \varphi(z, \cdot; x_0, \alpha(\xi)) \\ \vartheta'(z, \cdot; x_0, \alpha(\xi)) & \varphi'(z, \cdot; x_0, \alpha(\xi)) \end{bmatrix},\end{aligned}\quad (3.5)$$

in particular,

$$\Psi(z, x_0; x_0, \alpha(\xi)) = \begin{bmatrix} \cos(\xi) & -\sin(\xi) \\ \sin(\xi) & \cos(\xi) \end{bmatrix}, \quad x_0 \in [0, R].$$

In addition, set

$$\beta(\eta) = [\cos(\eta) \quad -\sin(\eta)] = \alpha(-\eta), \quad \eta \in S_{2\pi}. \quad (3.6)$$

As shown in [21, Section 2] and [22], one can then introduce for all $x_0, y_0 \in [0, R]$, $x_0 \neq y_0$, the *Weyl–Titchmarsh function*, $m(z; x_0, \alpha(\xi); y_0, \beta(\eta))$, given by

$$\begin{aligned}m(z; x_0, \alpha(\xi); y_0, \beta(\eta)) &= -[\beta(\eta)\Phi(z, y_0; x_0, \alpha(\xi))]^{-1}[\beta(\eta)\Theta(z, y_0; x_0, \alpha(\xi))] \\ &= -\frac{\cos(\eta)\vartheta(z, y_0; x_0, \alpha(\xi)) - \sin(\eta)\vartheta'(z, y_0; x_0, \alpha(\xi))}{\cos(\eta)\varphi(z, y_0; x_0, \alpha(\xi)) - \sin(\eta)\varphi'(z, y_0; x_0, \alpha(\xi))},\end{aligned}\quad (3.7)$$

$\xi, \eta \in S_{2\pi}, z \in \mathbb{C} \setminus \sigma(H_{\xi, \eta}).$

Then, with $\psi(z, \cdot; x_0, \alpha(\xi); y_0, \beta(\eta))$ defined by

$$\begin{aligned}\psi(z, \cdot; x_0, \alpha(\xi); y_0, \beta(\eta)) &= \vartheta(z, \cdot; x_0, \alpha(\xi)) + \varphi(z, \cdot; x_0, \alpha(\xi))m(z; x_0, \alpha(\xi); y_0, \beta(\eta)),\end{aligned}\quad (3.8)$$

$\varphi(z, \cdot; x_0, \alpha(\xi))$ and $\psi(z, \cdot; x_0, \alpha(\xi); y_0, \beta(\eta))$ are linearly independent solutions of (2.25) satisfying

$$\cos(\xi)\varphi(z, x_0; x_0, \alpha(\xi)) + \sin(\xi)\varphi'(z, x_0; x_0, \alpha(\xi)) = 0, \quad (3.9)$$

$$\cos(\eta)\psi(z, y_0; x_0, \alpha(\xi); y_0, \beta(\eta)) - \sin(\eta)\psi'(z, y_0; x_0, \alpha(\xi); y_0, \beta(\eta)) = 0. \quad (3.10)$$

The following result has been proved in [21, Lemma 2.10] in a self-adjoint context, but self-adjointness is of no relevance for the result (3.11) below. For the convenience of the reader we reproduce its proof here.

Lemma 3.1 ([21], Lemma 2.10). *Suppose $\xi, \eta \in S_{2\pi}$, $z \in \mathbb{C} \setminus (\sigma(H_{\xi, \eta}) \cup \sigma(H_{0, \eta}))$, let $\alpha(\xi)$ and $\beta(\eta)$ be defined as in (3.4) and (3.6), and suppose that $x_0, y_0 \in [0, R]$. Then, the following linear fractional transformation holds,*

$$m(z; x_0, \alpha(\xi); y_0, \beta(\eta)) = \frac{-\sin(\xi) + \cos(\xi)m(z; x_0, \alpha(0); y_0, \beta(\eta))}{\cos(\xi) + \sin(\xi)m(z; x_0, \alpha(0); y_0, \beta(\eta))}, \quad (3.11)$$

with m defined as in (3.7).

Proof. With $\psi(z, \cdot; x_0, \alpha(\xi); y_0, \beta(\eta))$ defined in (3.8), by (3.10) one obtains

$$\beta(\eta) \begin{bmatrix} \psi(z, y_0; x_0, \alpha(0); y_0, \beta(\eta)) \\ \psi'(z, y_0; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix} = \beta(\eta) \begin{bmatrix} \psi(z, y_0; x_0, \alpha(\xi); y_0, \beta(\eta)) \\ \psi'(z, y_0; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix} = 0, \quad (3.12)$$

and as a consequence, for some $c(z) \in \mathbb{C} \setminus \{0\}$,

$$\begin{bmatrix} \psi(z, y_0; x_0, \alpha(\xi); y_0, \beta(\eta)) \\ \psi'(z, y_0; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix} = c(z) \begin{bmatrix} \psi(z, y_0; x_0, \alpha(0); y_0, \beta(\eta)) \\ \psi'(z, y_0; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix}. \quad (3.13)$$

By the uniqueness of solutions for the initial value problem associated with (2.25),

$$\begin{bmatrix} \psi(z, x; x_0, \alpha(\xi); y_0, \beta(\eta)) \\ \psi'(z, x; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix} = c(z) \begin{bmatrix} \psi(z, x; x_0, \alpha(0); y_0, \beta(\eta)) \\ \psi'(z, x; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix}, \quad x \in [0, R]. \quad (3.14)$$

Next, one notes that

$$\begin{aligned} \begin{bmatrix} \psi(z, x_0; x_0, \alpha(\xi); y_0, \beta(\eta)) \\ \psi'(z, x_0; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix} &= \Psi(z, x_0; x_0, \alpha(\xi)) \begin{bmatrix} 1 \\ m(z; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix} \\ &= \begin{bmatrix} \cos(\xi) & -\sin(\xi) \\ \sin(\xi) & \cos(\xi) \end{bmatrix} \begin{bmatrix} 1 \\ m(z; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix}, \end{aligned} \quad (3.15)$$

and similarly that

$$\begin{aligned} \begin{bmatrix} \psi(z, x_0; x_0, \alpha(0); y_0, \beta(\eta)) \\ \psi'(z, x_0; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix} &= \Psi(z, x_0; x_0, \alpha(0)) \begin{bmatrix} 1 \\ m(z; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix} \\ &= \begin{bmatrix} 1 \\ m(z; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix}. \end{aligned} \quad (3.16)$$

By (3.14)–(3.16), one concludes that

$$\begin{bmatrix} 1 \\ m(z; x_0, \alpha(\xi); y_0, \beta(\eta)) \end{bmatrix} = c(z) \begin{bmatrix} \cos(\xi) & \sin(\xi) \\ -\sin(\xi) & \cos(\xi) \end{bmatrix} \begin{bmatrix} 1 \\ m(z; x_0, \alpha(0); y_0, \beta(\eta)) \end{bmatrix} \quad (3.17)$$

implying (3.11). \square

With $\alpha(\xi)$ and $\beta(\eta)$ defined in (3.4) and (3.6), and with $m(z; x_0, \alpha(\xi); y_0, \beta(\eta))$ defined in (3.7), for the next result we let

$$\alpha_0 = \alpha(\theta_0) = [\cos(\theta_0) \quad \sin(\theta_0)], \quad \beta_R = \beta(\theta_R) = [\cos(\theta_R) \quad -\sin(\theta_R)], \quad (3.18)$$

$$m_{+, \theta_0}(z, \theta_R) = m(z; 0, \alpha_0; R, \beta_R), \quad m_{-, \theta_R}(z, \theta_0) = m(z; R, \beta_R; 0, \alpha_0). \quad (3.19)$$

Theorem 3.2. *Let $\theta_0, \theta_R \in S_{2\pi}$ and $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$. Then,*

$$\Lambda_{\theta_0, \theta_R}(z) = \begin{bmatrix} m_{+, \theta_0}(z, \theta_R) & \Lambda_{\theta_0, \theta_R}(z)_{1,2} \\ \Lambda_{\theta_0, \theta_R}(z)_{2,1} & -m_{-, \theta_R}(z, \theta_0) \end{bmatrix}, \quad (3.20)$$

where

$$\begin{aligned}
\Lambda_{\theta_0, \theta_R}(z)_{1,2} &= \Lambda_{\theta_0, \theta_R}(z)_{2,1} \\
&= \frac{-\sin(\theta_0)u_{-, \theta_0}(z, 0) + \cos(\theta_0)u'_{-, \theta_0}(z, 0)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)} \\
&= \frac{-\sin(\theta_R)u_{+, \theta_R}(z, R) - \cos(\theta_R)u'_{+, \theta_R}(z, R)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)}.
\end{aligned} \tag{3.21}$$

Proof. We temporarily assume in addition that $z \notin (\sigma(H_{\theta_0, 0}) \cup \sigma(H_{0, \theta_R}))$. If $x_0 = 0, y_0 = R$, then with $\varphi(z, \cdot; 0, \alpha_0)$ defined in (3.5), and with $\psi(z, \cdot; 0, \alpha_0; R, \beta_R)$ defined in (3.8), one notes, by (3.9) and (3.10), that

$$\varphi(z, \cdot; 0, \alpha_0) = C_-(z)u_{-, \theta_0}(z, \cdot), \tag{3.22}$$

$$\psi(z, \cdot; 0, \alpha_0; R, \beta_R) = C_+(z)u_{+, \theta_R}(z, \cdot), \tag{3.23}$$

for some $C_{\pm}(z) \in \mathbb{C} \setminus \{0\}$, where $u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot)$ represents the basis for the solutions of (2.25) described in (2.33). With $m_{+, \theta_0}(z, \theta_R)$ defined in (3.19), one notes when $\theta_0 = 0$ that

$$m_{+, 0}(z, \theta_R) = \left. \frac{\psi'(z, 0; 0, \alpha_0; R, \beta_R)}{\psi(z, 0; 0, \alpha_0; R, \beta_R)} \right|_{\theta_0=0} = u'_{+, \theta_R}(z, 0); \tag{3.24}$$

hence by (3.11) that

$$\begin{aligned}
m_{+, \theta_0}(z, \theta_R) &= \frac{-\sin(\theta_0) + \cos(\theta_0)m_{+, 0}(z, \theta_R)}{\cos(\theta_0) + \sin(\theta_0)m_{+, 0}(z, \theta_R)} \\
&= \frac{-\sin(\theta_0) + \cos(\theta_0)u'_{+, \theta_R}(z, 0)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)}.
\end{aligned} \tag{3.25}$$

By Theorem 2.3, $\Lambda_{\theta_0, \theta_R}$ is invariant with respect to a change of basis for (2.25); thus the $(1, 1)$ -entry of $\Lambda_{\theta_0, \theta_R}$, provided in (3.3), equals $m_{+, \theta_0}(z, \theta_R)$.

Next, if $x_0 = R, y_0 = 0$, then with $\varphi(z, \cdot; R, \beta_R)$ defined in (3.5), and with $\psi(z, \cdot; R, \beta_R; 0, \alpha_0)$ defined in (3.8), one notes, by (3.9) and (3.10), that

$$\varphi(z, \cdot; R, \beta_R) = D_+(z)u_{+, \theta_R}(z, \cdot), \tag{3.26}$$

$$\psi(z, \cdot; R, \beta_R; 0, \alpha_0) = D_-(z)u_{-, \theta_0}(z, \cdot), \tag{3.27}$$

for some $D_{\pm}(z) \in \mathbb{C} \setminus \{0\}$, where $u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot)$ again denotes the basis for the solutions of (2.25) described in (2.33). With $m_{-, \theta_R}(z, \theta_0)$ defined in (3.19), one now obtains when $\theta_R = 0$ that,

$$m_{-, 0}(z, \theta_0) = \left. \frac{\psi'(z, R; R, \beta_R; 0, \alpha_0)}{\psi(z, R; R, \beta_R; 0, \alpha_0)} \right|_{\theta_R=0} = u'_{-, \theta_0}(z, R); \tag{3.28}$$

hence by (3.11) that

$$\begin{aligned} m_{-, \theta_R}(z, \theta_0) &= \frac{\sin(\theta_R) + \cos(\theta_R)m_{-,0}(z, \theta_0)}{\cos(\theta_R) - \sin(\theta_R)m_{-,0}(z, \theta_0)} \\ &= \frac{\sin(\theta_R) + \cos(\theta_R)u'_{-, \theta_0}(z, R)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)}. \end{aligned} \quad (3.29)$$

Again, by the invariance of $\Lambda_{\theta_0, \theta_R}$ with respect to a change of basis for solutions of (2.25), the $(2, 2)$ -entry of $\Lambda_{\theta_0, \theta_R}$, provided in (3.3), is given by $-m_{-, \theta_R}(z, \theta_0)$.

To see that the off-diagonal elements of (3.20) are equal, we introduce

$$W(z) = W(u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot)), \quad \theta_0, \theta_R \in S_{2\pi}, \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad (3.30)$$

where $u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot)$, is the basis for the solutions of (2.25) as described in (2.33). A straightforward computation, then yields

$$\begin{aligned} &[-\sin(\theta_0)u_{-, \theta_0}(z, 0) + \cos(\theta_0)u'_{-, \theta_0}(z, 0)][\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)] \\ &= W(z) \\ &= [\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)][-\sin(\theta_R)u_{+, \theta_R}(z, R) - \cos(\theta_R)u'_{+, \theta_R}(z, R)]. \end{aligned} \quad (3.31)$$

Indeed, the first equality in (3.31) follows by inserting the second term in (2.34) into (2.37); similarly, the second equality in (3.31) follows by inserting the second term in (2.35) into (2.38). Equation (3.31) immediately yields (3.21).

Finally, a meromorphic continuation with respect to z then removes the additional assumption $z \notin (\sigma(H_{\theta_0, 0}) \cup \sigma(H_{0, \theta_R}))$, completing the proof. \square

Remark 3.3. Assume the special self-adjoint case where V is real-valued and $\theta_0, \theta_R \in [0, 2\pi)$. By Marchenko's fundamental uniqueness result [59, Ch. 2], one observes the inverse spectral theory fact that each of the two diagonal terms of $\Lambda_{\theta_0, \theta_R}(\cdot)$ already uniquely determines the potential coefficient $V(x)$ for a.e. $x \in [0, R]$. In addition, the known leading asymptotic behavior of $m_{+, \theta_0}(z, \theta_R)$ as $z \rightarrow i\infty$,

$$\begin{aligned} m_{+, \theta_0}(z, \theta_R) &\xrightarrow{z \rightarrow i\infty} \cot(\theta_0) + o(1), \quad \theta_0 \in [0, 2\pi) \setminus \{0, \pi\}, \\ m_{+, 0}(z, \theta_R) &\xrightarrow{z \rightarrow i\infty} iz^{1/2} + o(z^{1/2}), \quad \theta_0 \in \{0, \pi\}, \end{aligned} \quad (3.32)$$

determines the boundary condition parameter θ_0 , and similarly, that of $m_{-, \theta_R}(z, \theta_0)$ determines θ_R . These facts are obviously not shared by the usual 2×2 matrix-valued Weyl–Titchmarsh function to be briefly discussed in our final Section 6. and hence demonstrates a stark contrast between these two matrix-valued functions, even though, both are Herglotz functions. The Herglotz property is well-known to be an intrinsic property of the 2×2 matrix-valued Weyl–Titchmarsh function (cf., e.g., [23, Sect. 9.5]), [57, Sect. II.8], [68, Sect. 6.5], [80, Ch. III]) and will also be verified for $\Lambda_{\theta_0, \theta_R}(\cdot)$ as a special case of the principal result in Theorem 4.6 (cf. (4.77)).

Next we turn to the asymptotic behavior of $\Lambda_{\theta_0, \theta_R}(z)$ as $|z| \rightarrow \infty$, $\text{Im}(z^{1/2}) > 0$:

Lemma 3.4. *Suppose that $V \in L^1((0, R); dx)$ and let $\theta_0, \theta_R \in S_{2\pi}$. Then,*

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}(z) &\stackrel{=}{=} \lim_{\substack{|z| \rightarrow \infty \\ \text{Im}(z^{1/2}) > 0}} \begin{bmatrix} \cot(\theta_0) & \frac{2iz^{-1/2}e^{iz^{1/2}R}}{\sin(\theta_0)\sin(\theta_R)} \\ \frac{2iz^{-1/2}e^{iz^{1/2}R}}{\sin(\theta_0)\sin(\theta_R)} & \cot(\theta_0) \end{bmatrix} \\ &+ \begin{bmatrix} O(|z|^{-1/2}) & O(|z|^{-1}e^{-\text{Im}(z^{1/2})R}) \\ O(|z|^{-1}e^{-\text{Im}(z^{1/2})R}) & O(|z|^{-1/2}) \end{bmatrix}, \quad \theta_0 \neq 0, \theta_R \neq 0, \end{aligned} \quad (3.33)$$

$$\begin{aligned} \Lambda_{0, \theta_R}(z) &\stackrel{=}{=} \lim_{\substack{|z| \rightarrow \infty \\ \text{Im}(z^{1/2}) > 0}} \begin{bmatrix} iz^{1/2} & -\frac{2e^{iz^{1/2}R}}{\sin(\theta_R)} \\ -\frac{2e^{iz^{1/2}R}}{\sin(\theta_R)} & \cot(\theta_R) \end{bmatrix} \\ &+ \begin{bmatrix} O(1) & O(|z|^{-1/2}e^{-\text{Im}(z^{1/2})R}) \\ O(|z|^{-1/2}e^{-\text{Im}(z^{1/2})R}) & O(|z|^{-1/2}) \end{bmatrix}, \quad \theta_0 = 0, \theta_R \neq 0, \end{aligned} \quad (3.34)$$

$$\begin{aligned} \Lambda_{\theta_0, 0}(z) &\stackrel{=}{=} \lim_{\substack{|z| \rightarrow \infty \\ \text{Im}(z^{1/2}) > 0}} \begin{bmatrix} \cot(\theta_0) & -\frac{2e^{iz^{1/2}R}}{\sin(\theta_0)} \\ -\frac{2e^{iz^{1/2}R}}{\sin(\theta_0)} & iz^{1/2} \end{bmatrix} \\ &+ \begin{bmatrix} O(|z|^{-1/2}) & O(|z|^{-1/2}e^{-\text{Im}(z^{1/2})R}) \\ O(|z|^{-1/2}e^{-\text{Im}(z^{1/2})R}) & O(1) \end{bmatrix}, \quad \theta_0 \neq 0, \theta_R = 0, \end{aligned} \quad (3.35)$$

$$\begin{aligned} \Lambda_{0, 0}(z) &\stackrel{=}{=} \lim_{\substack{|z| \rightarrow \infty \\ \text{Im}(z^{1/2}) > 0}} \begin{bmatrix} iz^{1/2} & -2iz^{1/2}e^{iz^{1/2}R} \\ -2iz^{1/2}e^{iz^{1/2}R} & iz^{1/2} \end{bmatrix} \\ &+ \begin{bmatrix} O(1) & O(e^{-\text{Im}(z^{1/2})R}) \\ O(e^{-\text{Im}(z^{1/2})R}) & O(1) \end{bmatrix}, \quad \theta_0 = 0, \theta_R = 0. \end{aligned} \quad (3.36)$$

Proof. This follows from an elementary computation upon inserting (2.55) (resp., (2.56)) into (3.25) (resp., (3.29)) and using the asymptotic expansions (2.18). \square

For the principal result of this section, an explicit formula for $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ in terms of the resolvent $(H_{\theta_0, \theta_R} - zI)^{-1}$ of H_{θ_0, θ_R} and the boundary traces $\gamma_{\theta'_0, \theta'_R}$, we recall the Green's function associated with the operator H_{θ_0, θ_R} in (2.5),

$$\begin{aligned} G_{\theta_0, \theta_R}(z, x, x') &= (H_{\theta_0, \theta_R} - zI)^{-1}(x, x') \\ &= \frac{1}{W(u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot))} \begin{cases} u_{-, \theta_0}(z, x')u_{+, \theta_R}(z, x), & 0 \leq x' \leq x, \\ u_{-, \theta_0}(z, x)u_{+, \theta_R}(z, x'), & 0 \leq x \leq x', \end{cases} \\ &\quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad x, x' \in [0, R]. \end{aligned} \quad (3.37)$$

Here $u_{+, \theta_R}(z, \cdot)$, $u_{-, \theta_0}(z, \cdot)$ is a basis for solutions of (2.25) as described in (2.33), and we abbreviate the identity operator in $L^2((0, R); dx)$ by writing $I = I_{L^2((0, R); dx)}$. Thus, one obtains

$$\begin{aligned} ((H_{\theta_0, \theta_R} - zI)^{-1}g)(x) &= \int_0^R dx' G_{\theta_0, \theta_R}(z, x, x')g(x'), \\ g &\in L^2((0, R); dx), \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad x \in (0, R). \end{aligned} \quad (3.38)$$

Equation (3.37) is a crucial input in the proof of the following fundamental result:

Theorem 3.5. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, let H_{θ_0, θ_R} be defined as in (2.5), suppose that $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$, and let $S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ be defined according to (2.43). Then*

$$\Lambda_{\theta'_0, \theta'_R}^{\theta_0, \theta_R}(z) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} = \gamma_{\theta'_0, \theta'_R} [\overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*. \quad (3.39)$$

In particular, with $\theta'_0 = (\theta_0 + \pi/2) \bmod(2\pi)$, $\theta'_R = (\theta_R + \pi/2) \bmod(2\pi)$, one obtains

$$\Lambda_{\theta_0, \theta_R}(z) = \widehat{\gamma}_{\theta_0, \theta_R} [\widehat{\overline{\gamma_{\theta_0, \theta_R}}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*, \quad (3.40)$$

where

$$\widehat{\gamma}_{\theta_0, \theta_R} = \gamma_{(\theta_0 + \frac{\pi}{2}) \bmod(2\pi), (\theta_R + \frac{\pi}{2}) \bmod(2\pi)}. \quad (3.41)$$

As a consequence,

$$\Lambda_{(\theta_0 + \pi) \bmod(2\pi), (\theta_R + \pi) \bmod(2\pi)}(z) = \Lambda_{\theta_0, \theta_R}(z), \quad \theta_0, \theta_R \in S_{2\pi}. \quad (3.42)$$

Proof. We start by noting that the Green's function $G_{\theta_0, \theta_R}^*(z, x, x')$ of $(H_{\theta_0, \theta_R})^*$ is given by

$$\begin{aligned} G_{\theta_0, \theta_R}^*(z, x, x') &= ((H_{\theta_0, \theta_R})^* - zI)^{-1}(x, x') \\ &= \frac{1}{W^*(z)} \begin{cases} \overline{u_{-, \theta_0}(\bar{z}, x')} \overline{u_{+, \theta_R}(\bar{z}, x)}, & 0 \leq x' \leq x, \\ \overline{u_{-, \theta_0}(\bar{z}, x)} \overline{u_{+, \theta_R}(\bar{z}, x')}, & 0 \leq x \leq x', \end{cases} \\ &\quad z \in \mathbb{C} \setminus \sigma((H_{\theta_0, \theta_R})^*), \quad x, x' \in [0, R], \end{aligned}$$

where $u_{+, \theta_R}(z, \cdot)$, $u_{-, \theta_0}(z, \cdot)$ is a basis for solutions of (2.25) as described in (2.33), and $W^*(z)$ is defined by

$$W^*(z) = W(\overline{u_{+, \theta_R}(\bar{z}, \cdot)}, \overline{u_{-, \theta_0}(\bar{z}, \cdot)}), \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}). \quad (3.43)$$

In particular, we note the fact that

$$\overline{W^*(\bar{z})} = W(z), \quad (3.44)$$

with $W(z)$ the Wronskian defined in (3.30). Thus, one obtains

$$\begin{aligned} (((H_{\theta_0, \theta_R})^* - zI)^{-1}g)(x) &= \int_0^R dx' G_{\theta_0, \theta_R}^*(z, x, x')g(x'), \\ g &\in L^2((0, R); dx), \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad x \in (0, R). \end{aligned} \quad (3.45)$$

Next, for $g \in L^2((0, R); dx)$ let

$$\begin{aligned}\widehat{g}(\bar{z}, x) &= (((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}g)(x) \\ &= \frac{1}{W(z)} \left(\overline{u_{+, \theta_R}(z, x)} \int_0^x dx' \overline{u_{-, \theta_0}(z, x')} g(x') \right. \\ &\quad \left. + \overline{u_{-, \theta_0}(z, x)} \int_x^R dx' \overline{u_{+, \theta_R}(z, x')} g(x') \right).\end{aligned}\tag{3.46}$$

Then one notes that

$$\begin{aligned}\widehat{g}'(\bar{z}, x) &= \frac{1}{W(z)} \left(\overline{u'_{+, \theta_R}(z, x)} \int_0^x dx' \overline{u_{-, \theta_0}(z, x')} g(x') \right. \\ &\quad \left. + \overline{u'_{-, \theta_0}(z, x)} \int_x^R dx' \overline{u_{+, \theta_R}(z, x')} g(x') \right),\end{aligned}\tag{3.47}$$

and, as a consequence, that

$$\widehat{g}(\bar{z}, 0) = \frac{1}{W(z)} \overline{u_{-, \theta_0}(z, 0)} \int_0^R dx' \overline{u_{+, \theta_R}(z, x')} g(x'),\tag{3.48}$$

$$\widehat{g}'(\bar{z}, 0) = \frac{1}{W(z)} \overline{u'_{-, \theta_0}(z, 0)} \int_0^R dx' \overline{u_{+, \theta_R}(z, x')} g(x'),\tag{3.49}$$

$$\widehat{g}(\bar{z}, R) = \frac{1}{W(z)} \overline{u_{+, \theta_R}(z, R)} \int_0^R dx' \overline{u_{-, \theta_0}(z, x')} g(x'),\tag{3.50}$$

$$\widehat{g}'(\bar{z}, R) = \frac{1}{W(z)} \overline{u'_{+, \theta_R}(z, R)} \int_0^R dx' \overline{u_{-, \theta_0}(z, x')} g(x').\tag{3.51}$$

In turn, this permits one to compute

$$\begin{aligned}\gamma_{\theta'_0, \theta'_R} \overline{((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}g} &= \gamma_{\theta'_0, \theta'_R} \widehat{g}(\bar{z}, \cdot) \\ &= \begin{bmatrix} \cos(\overline{\theta'_0}) \widehat{g}(\bar{z}, 0) + \sin(\overline{\theta'_0}) \widehat{g}'(\bar{z}, 0) \\ \cos(\overline{\theta'_R}) \widehat{g}(\bar{z}, R) - \sin(\overline{\theta'_R}) \widehat{g}'(\bar{z}, R) \end{bmatrix}.\end{aligned}\tag{3.52}$$

Using (3.48)–(3.52) one infers, with $[a_0 \ a_R]^\top \in \mathbb{C}^2$ and $(\cdot, \cdot)_{\mathbb{C}^2}$ denoting the scalar product in \mathbb{C}^2 , that

$$\begin{aligned}&(\gamma_{\theta'_0, \theta'_R} \overline{((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}g}, [a_0 \ a_R]^\top)_{\mathbb{C}^2} \\ &= \left(\begin{bmatrix} \cos(\overline{\theta'_0}) \widehat{g}(\bar{z}, 0) + \sin(\overline{\theta'_0}) \widehat{g}'(\bar{z}, 0) \\ \cos(\overline{\theta'_R}) \widehat{g}(\bar{z}, R) - \sin(\overline{\theta'_R}) \widehat{g}'(\bar{z}, R) \end{bmatrix}, \begin{bmatrix} a_0 \\ a_R \end{bmatrix} \right)_{\mathbb{C}^2} \\ &= \frac{1}{W(z)} \int_0^R dx' \overline{g(x')} (\cos(\theta'_0) a_0 \overline{u_{-, \theta_0}(z, 0)} u_{+, \theta_R}(z, x') \\ &\quad + \sin(\theta'_0) a_0 \overline{u'_{-, \theta_0}(z, 0)} u_{+, \theta_R}(z, x') + \cos(\theta'_R) a_R \overline{u_{+, \theta_R}(z, R)} u_{-, \theta_0}(z, x') \\ &\quad - \sin(\theta'_R) a_R \overline{u'_{+, \theta_R}(z, R)} u_{-, \theta_0}(z, x')).\end{aligned}\tag{3.53}$$

Hence one concludes that

$$\begin{aligned} & \left([\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [a_0 \ a_R]^\top \right) (x) \\ &= \frac{1}{W(z)} \left([\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] a_0 u_{+, \theta_R}(z, x) \right. \\ & \quad \left. + [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] a_R u_{-, \theta_0}(z, x) \right). \end{aligned} \quad (3.54)$$

Consequently, writing

$$\begin{aligned} & \gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [a_0 \ a_R]^\top \\ &= \begin{bmatrix} (\gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [a_0 \ a_R]^\top)_1 \\ (\gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [a_0 \ a_R]^\top)_2 \end{bmatrix}, \end{aligned} \quad (3.55)$$

it follows that

$$\begin{aligned} & (\gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [a_0 \ a_R]^\top)_1 \\ &= \frac{1}{W(z)} \left(\cos(\theta'_0) [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] a_0 \right. \\ & \quad + \cos(\theta'_0) [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] a_R u_{-, \theta_0}(z, 0) \\ & \quad + \sin(\theta'_0) [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] a_0 u'_{+, \theta_R}(z, 0) \\ & \quad \left. + \sin(\theta'_0) [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] a_R u'_{-, \theta_0}(z, 0) \right), \end{aligned} \quad (3.56)$$

and

$$\begin{aligned} & (\gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [a_0 \ a_R]^\top)_2 \\ &= \frac{1}{W(z)} \left(\cos(\theta'_R) [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] a_0 u_{+, \theta_R}(z, R) \right. \\ & \quad + \cos(\theta'_R) [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] a_R \\ & \quad - \sin(\theta'_R) [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] a_0 u'_{+, \theta_R}(z, R) \\ & \quad \left. - \sin(\theta'_R) [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] a_R u'_{-, \theta_0}(z, R) \right). \end{aligned} \quad (3.57)$$

Hence, writing

$$\begin{aligned} & \gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* \\ &= \left[(\gamma_{\theta'_0, \theta'_R} [\gamma_{\theta'_0, \theta'_R} \overline{\gamma_{\theta'_0, \theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{j,k} \right]_{j,k=1,2}, \end{aligned} \quad (3.58)$$

one obtains

$$\begin{aligned}
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{1,1} \\
&= \frac{1}{W(z)} [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] \\
&\quad \times [\cos(\theta'_0) + \sin(\theta'_0) u'_{+, \theta_R}(z, 0)], \\
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{1,2} \\
&= \frac{1}{W(z)} [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] \\
&\quad \times [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)], \\
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{2,1} \\
&= \frac{1}{W(z)} [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] \\
&\quad \times [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)], \\
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{2,2} \\
&= \frac{1}{W(z)} [\cos(\theta'_R) u_{+, \theta_R}(z, R) - \sin(\theta'_R) u'_{+, \theta_R}(z, R)] \\
&\quad \times [\cos(\theta'_R) - \sin(\theta'_R) u'_{-, \theta_0}(z, R)].
\end{aligned} \tag{3.59}$$

Employing (2.37) and (2.38) one finally concludes from (2.51) that the expressions in (3.59) are equivalent to the following ones

$$\begin{aligned}
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{1,1} = \sin(\theta'_0 - \theta_0) \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{1,1}, \\
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{1,2} = \sin(\theta'_R - \theta_R) \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{1,2}, \\
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{2,1} = \sin(\theta'_0 - \theta_0) \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{2,1}, \\
& (\gamma_{\theta'_0, \theta'_R} [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^*)_{2,2} = \sin(\theta'_R - \theta_R) \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right)_{2,2},
\end{aligned} \tag{3.60}$$

proving (3.39).

Finally, equation (3.42) is a consequence of (2.6) and (3.40). \square

Remark 3.6. A formula of the type (3.40) for Dirichlet-to-Neumann maps associated with multi-dimensional Schrödinger operators was published by Amrein and Pearson [4] in 2004. It has recently been extended in various directions in [35], [36], [38], [39]. Formula (3.40) for the Dirichlet-to-Neumann map $\Lambda_{\pi/2, \pi/2}$ in the special case $V = 0$ has also been derived by Posilicano [72, Example 5.1].

Remark 3.7. While it is tempting to view $\gamma_{\theta_0, \theta_R}$ as an unbounded but densely defined operator in $L^2((0, R); dx)$ whose domain contains the space $C_0^\infty((0, R))$, one should note, in this case, that

its adjoint $\gamma_{\theta_0, \theta_R}^*$ is not densely defined. Indeed, the adjoint $\gamma_{\theta_0, \theta_R}^*$ of $\gamma_{\theta_0, \theta_R}$ would have to be an unbounded operator from \mathbb{C}^2 to $L^2((0, R); dx)$ such that

$$(\gamma_{\theta_0, \theta_R} f, g)_{\mathbb{C}^2} = (f, \gamma_{\theta_0, \theta_R}^* g)_{L^2((0, R); dx)} \text{ for all } f \in \text{dom}(\gamma_{\theta_0, \theta_R}), g \in \text{dom}(\gamma_{\theta_0, \theta_R}^*). \quad (3.61)$$

In particular, choosing $f \in C_0^\infty((0, R))$, in which case $\gamma_{\theta_0, \theta_R} f = 0$, one concludes that

$$(f, \gamma_{\theta_0, \theta_R}^* g)_{L^2((0, R); dx)} = 0 \text{ for all } f \in C_0^\infty((0, R)). \quad (3.62)$$

Thus, one obtains $\gamma_{\theta_0, \theta_R}^* g = 0$ for all $g \in \text{dom}(\gamma_{\theta_0, \theta_R}^*)$. Since obviously $\gamma_{\theta_0, \theta_R} \neq 0$, (3.61) implies $\text{dom}(\gamma_{\theta_0, \theta_R}^*) = \{0\}$ and hence $\gamma_{\theta_0, \theta_R}$ is not a closable linear operator in $L^2((0, R); dx)$. This is the reason for our careful choice of notation in (3.39) and (3.40).

We conclude this section by providing an explicit example in the special case $V = 0$ a.e. on $[0, R]$. For notational purposes we add the superscript (0) to the corresponding quantities below:

Example 3.8. Let $V = 0$ a.e. on $[0, R]$, $x_0 \in (0, R)$, $\theta_0, \theta_R, \theta \in S_{2\pi}$, $x, x', s \in [0, R]$, and let

$$f(z, s, \alpha, \beta) = f(z, s, \beta, \alpha) \quad (3.63)$$

$$\begin{aligned} &= z \sin(\alpha) \sin(\beta) \sin(\sqrt{z}s) + \sqrt{z} \sin(\alpha + \beta) \cos(\sqrt{z}s) - \cos(\alpha) \cos(\beta) \sin(\sqrt{z}s), \\ g(z, s, \alpha, \beta) &= f(z, s, \alpha + \pi/2, \beta) \\ &= z \cos(\alpha) \sin(\beta) \sin(\sqrt{z}s) + \sqrt{z} \cos(\alpha + \beta) \cos(\sqrt{z}s) + \sin(\alpha) \cos(\beta) \sin(\sqrt{z}s). \end{aligned} \quad (3.64)$$

Then,

$$u^{(0)}(z, x, (\theta_0, c_0), (\theta_R, c_R)) = \frac{c_0 f(z, R - x, 0, \theta_R) + c_R f(z, x, \theta_0, 0)}{f(z, R, \theta_0, \theta_R)}, \quad (3.65)$$

$$u_{+, \theta_R}^{(0)}(z, x) = u(z, \cdot; (0, 1), (\theta_R, 0)) = \frac{f(z, R - x, 0, \theta_R)}{f(z, R, 0, \theta_R)}, \quad (3.66)$$

$$u_{+, \theta_R}^{(0)'}(z, x) = \frac{f'(z, R - x, 0, \theta_R)}{f(z, R, 0, \theta_R)} = \frac{f(z, R - x, \pi/2, \theta_R)}{f(z, R, 0, \theta_R)}, \quad (3.67)$$

$$u_{-, \theta_0}^{(0)}(z, x) = u(z, \cdot; (\theta_0, 0), (0, 1)) = \frac{f(z, x, \theta_0, 0)}{f(z, R, \theta_0, 0)}, \quad (3.68)$$

$$u_{-, \theta_0}^{(0)'}(z, x) = \frac{f'(z, x, \theta_0, 0)}{f(z, R, \theta_0, 0)} = -\frac{f(z, x, \theta_0, \pi/2)}{f(z, R, \theta_0, 0)}, \quad (3.69)$$

$$W(u_{+, \theta_R}^{(0)}(z, \cdot), u_{-, \theta_0}^{(0)}(z, \cdot)) = \frac{-\sqrt{z} f(z, R, \theta_0, \theta_R)}{f(z, R, \theta_0, 0) f(z, R, 0, \theta_R)}, \quad (3.70)$$

$$m_{+, \theta_0}^{(0)}(z, \theta_R) = \frac{g(z, R, \theta_0, \theta_R)}{f(z, R, \theta_0, \theta_R)}, \quad (3.71)$$

$$m_{-, \theta_R}^{(0)}(z, \theta_0) = -\frac{g(z, R, \theta_R, \theta_0)}{f(z, R, \theta_R, \theta_0)}, \quad (3.72)$$

$$\Lambda_{\theta_0, \theta_R}^{(0)}(z) = \begin{pmatrix} \frac{g(z, R, \theta_0, \theta_R)}{f(z, R, \theta_0, \theta_R)} & \frac{-\sqrt{z}}{f(z, R, \theta_0, \theta_R)} \\ \frac{-\sqrt{z}}{f(z, R, \theta_0, \theta_R)} & \frac{g(z, R, \theta_R, \theta_0)}{f(z, R, \theta_0, \theta_R)} \end{pmatrix},$$

$$m_{+,0}^{(0)}(z, x_0, \theta_R) = \frac{u_{+,\theta_R}^{(0)'}(z, x_0)}{u_{+,\theta_R}^{(0)}(z, x_0)} = \frac{g(z, R - x_0, 0, \theta_R)}{f(z, R - x_0, 0, \theta_R)}, \quad (3.73)$$

$$m_{-,0}^{(0)}(z, x_0, \theta_0) = \frac{u_{-,\theta_0}^{(0)'}(z, x_0)}{u_{-,\theta_0}^{(0)}(z, x_0)} = -\frac{g(z, x_0, 0, \theta_0)}{f(z, x_0, 0, \theta_0)}, \quad (3.74)$$

$$\begin{aligned} G_{\theta_0, \theta_R}^{(0)}(z, x, x') &= (H_{\theta_0, \theta_R}^{(0)} - zI)^{-1}(x, x') \\ &= \frac{1}{W(u_{+,\theta_R}^{(0)}(z, \cdot), u_{-,\theta_0}^{(0)}(z, \cdot))} \begin{cases} u_{-,\theta_0}^{(0)}(z, x')u_{+,\theta_R}^{(0)}(z, x), & 0 \leq x' \leq x, \\ u_{-,\theta_0}^{(0)}(z, x)u_{+,\theta_R}^{(0)}(z, x'), & 0 \leq x \leq x', \end{cases} \\ &= \begin{cases} \frac{f(z, x', \theta_0, 0)f(z, R - x, 0, \theta_R)}{-\sqrt{z}f(z, R, \theta_0, \theta_R)}, & 0 \leq x' \leq x, \\ \frac{f(z, x, \theta_0, 0)f(z, R - x', 0, \theta_R)}{-\sqrt{z}f(z, R, \theta_0, \theta_R)}, & 0 \leq x \leq x', \end{cases} \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}^{(0)}). \end{aligned} \quad (3.75)$$

The special case of the Neumann-to-Dirichlet map, $\Lambda_{\pi/2, \pi/2}^{(0)}$ in (3.73), was computed in [27, Example 4.1], and more recently, the special case of the Dirichlet-to-Neumann map, $\Lambda_{0,0}^{(0)}$ in (3.73), was computed in [72, Example 5.1].

4. Linear Fractional Transformations and the Herglotz Property in the Self-Adjoint Case

The principal purpose of this section is to prove that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ and $\Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)$ satisfy a linear fractional transformation. As a consequence we will show that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ is a 2×2 matrix-valued Herglotz function in the special self-adjoint case where V and $\theta_0, \theta_R, \theta'_0, \theta'_R$ are real-valued.

In the following we denote by $\mathbb{C}^{n \times n}$, $n \in \mathbb{N}$, the set of $n \times n$ matrices with complex-valued entries, and by I_n the identity matrix in \mathbb{C}^n .

Let $A = [A_{j,k}]_{1 \leq j,k \leq 2} \in \mathbb{C}^{4 \times 4}$, with $A_{j,k} \in \mathbb{C}^{2 \times 2}$, $1 \leq j, k \leq 2$, and $L \in \mathbb{C}^{2 \times 2}$, chosen such that $\ker(A_{1,1} + A_{1,2}L) = \{0\}$; that is, $(A_{1,1} + A_{1,2}L)$ is invertible in \mathbb{C}^2 . Define for such A (cf., e.g., [52]),

$$M_A(L) = (A_{2,1} + A_{2,2}L)(A_{1,1} + A_{1,2}L)^{-1}, \quad (4.1)$$

and observe that

$$M_{I_A}(L) = L, \quad (4.2)$$

$$M_{AB}(L) = M_A(M_B(L)), \quad (4.3)$$

$$M_{A^{-1}}(M_A(L)) = L = M_A(M_{A^{-1}}(L)), \quad A \text{ invertible}, \quad (4.4)$$

$$M_A(L) = M_{AB^{-1}}(M_B(L)), \quad (4.5)$$

whenever the right-hand sides (and hence the left-hand sides) in (4.3)–(4.5) exist.

Theorem 4.1. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R, \delta_0, \delta_R, \delta'_0, \delta'_R \in S_{2\pi}$, $\delta'_0 - \delta_0 \neq 0 \pmod{\pi}$, $\delta'_R - \delta_R \neq 0 \pmod{\pi}$, and that $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\delta_0, \delta_R}))$. Then, with S_{θ_0, θ_R} defined as in (2.43),*

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) &= (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} [S_{\delta'_0 - \theta'_0, \delta'_R - \theta'_R} + S_{\theta'_0 - \delta_0, \theta'_R - \delta_R} \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)] \\ &\quad \times [S_{\delta'_0 - \theta_0, \delta'_R - \theta_R} + S_{\theta_0 - \delta_0, \theta_R - \delta_R} \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)]^{-1} S_{\delta'_0 - \delta_0, \delta'_R - \delta_R}. \end{aligned} \quad (4.6)$$

Proof. Assume $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$, and let $\psi_j(z, \cdot)$, $j = 1, 2$, denote a basis for the solutions of (2.25). Then, with C_{θ_0, θ_R} and S_{θ_0, θ_R} as defined in (2.43), equation (2.42) yields

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) &\left(C_{\theta_0, \theta_R} \begin{bmatrix} \psi_1(z, 0) & \psi_2(z, 0) \\ \psi_1(z, R) & \psi_2(z, R) \end{bmatrix} + S_{\theta_0, \theta_R} \begin{bmatrix} \psi'_1(z, 0) & \psi'_2(z, 0) \\ -\psi'_1(z, R) & -\psi'_2(z, R) \end{bmatrix} \right) \\ &= \left(C_{\theta'_0, \theta'_R} \begin{bmatrix} \psi_1(z, 0) & \psi_2(z, 0) \\ \psi_1(z, R) & \psi_2(z, R) \end{bmatrix} + S_{\theta'_0, \theta'_R} \begin{bmatrix} \psi'_1(z, 0) & \psi'_2(z, 0) \\ -\psi'_1(z, R) & -\psi'_2(z, R) \end{bmatrix} \right). \end{aligned} \quad (4.7)$$

From Remark 2.6, recall, for $z \in \mathbb{C} \setminus \sigma(H_{0,0})$, that $\Lambda_{D,N}(z) = \Lambda_{0,0}^{\frac{\pi}{2}, \frac{\pi}{2}}(z) = \Lambda_{0,0}(z)$ and note that (2.45) then yields

$$\Lambda_{0,0}(z) = \begin{bmatrix} \psi'_1(z, 0) & \psi'_2(z, 0) \\ -\psi'_1(z, R) & -\psi'_2(z, R) \end{bmatrix} \begin{bmatrix} \psi_1(z, 0) & \psi_2(z, 0) \\ \psi_1(z, R) & \psi_2(z, R) \end{bmatrix}^{-1}. \quad (4.8)$$

Then, with C_{θ_0, θ_R} defined as in (2.43), and with $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{0,0}))$, (4.7) can be written as

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) = [C_{\theta'_0, \theta'_R} + S_{\theta'_0, \theta'_R} \Lambda_{0,0}(z)] [C_{\theta_0, \theta_R} + S_{\theta_0, \theta_R} \Lambda_{0,0}(z)]^{-1}. \quad (4.9)$$

Next, assume that $\theta_0, \theta_R, \theta'_0, \theta'_R, \delta_0, \delta_R, \delta'_0, \delta'_R \in S_{2\pi}$, and let $A, B \in \mathbb{C}^{4 \times 4}$ be defined by

$$A = \begin{bmatrix} C_{\theta_0, \theta_R} & S_{\theta_0, \theta_R} \\ C_{\theta'_0, \theta'_R} & S_{\theta'_0, \theta'_R} \end{bmatrix}, \quad B = \begin{bmatrix} C_{\delta_0, \delta_R} & S_{\delta_0, \delta_R} \\ C_{\delta'_0, \delta'_R} & S_{\delta'_0, \delta'_R} \end{bmatrix}. \quad (4.10)$$

Then, by (4.1), and (4.9),

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) = M_A(\Lambda_{0,0}(z)), \quad \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z) = M_B(\Lambda_{0,0}(z)) \quad (4.11)$$

for $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\delta_0, \delta_R}) \cup \sigma(H_{0,0}))$. If additionally, one assumes that $\delta'_0 - \delta_0 \neq 0 \pmod{\pi}$, $\delta'_R - \delta_R \neq 0 \pmod{\pi}$, then

$$AB^{-1} = \begin{bmatrix} (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\delta'_0 - \theta_0, \delta'_R - \theta_R} & (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\theta_0 - \delta_0, \theta_R - \delta_R} \\ (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\delta'_0 - \theta'_0, \delta'_R - \theta'_R} & (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\theta'_0 - \delta_0, \theta'_R - \delta_R} \end{bmatrix}, \quad (4.12)$$

and (4.6) then follows from (4.1), (4.3), (4.11) and (4.12), given that

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) = M_A(\Lambda_{0,0}(z)) = M_{AB^{-1}}(M_B(\Lambda_{0,0}(z))) = M_{AB^{-1}}(\Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)). \quad (4.13)$$

By meromorphic continuation, (4.6) holds for $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\delta_0, \delta_R}))$. \square

If H_{θ_0, θ_R} and H_{δ_0, δ_R} are self-adjoint, then (4.6) holds for $z \in \mathbb{C} \setminus \mathbb{R}$.

Remark 4.2. *In the special case of (4.6) which relates two generalized Dirichlet-to-Neumann maps one obtains*

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}(z) &= \left[-S_{\theta_0 - \delta_0, \theta_R - \delta_R} + C_{\theta_0 - \delta_0, \theta_R - \delta_R} \Lambda_{\delta_0, \delta_R}(z) \right] \\ &\quad \times \left[C_{\theta_0 - \delta_0, \theta_R - \delta_R} + S_{\theta_0 - \delta_0, \theta_R - \delta_R} \Lambda_{\delta_0, \delta_R}(z) \right]^{-1}, \\ &\quad \theta_0, \theta_R, \delta_0, \delta_R \in S_{2\pi}, \quad z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\delta_0, \delta_R})). \end{aligned} \quad (4.14)$$

The following reformulation of (4.6) (motivated by the form of (3.39) in Theorem 3.5) will be crucial in the proof of Theorem 4.6.

Corollary 4.3. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R, \delta_0, \delta_R, \delta'_0, \delta'_R \in S_{2\pi}$, $\theta'_0 - \theta_0 \neq 0 \pmod{\pi}$, $\theta'_R - \theta_R \neq 0 \pmod{\pi}$, $\delta'_0 - \delta_0 \neq 0 \pmod{\pi}$, $\delta'_R - \delta_R \neq 0 \pmod{\pi}$, and that $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\delta_0, \delta_R}))$. Then*

$$\begin{aligned} &\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} \\ &= \left[S_{\delta'_0 - \theta'_0, \delta'_R - \theta'_R} + (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\theta'_0 - \delta_0, \theta'_R - \delta_R} \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z) S_{\delta'_0 - \delta_0, \delta'_R - \delta_R} \right] \\ &\quad \times \left[(S_{\theta'_0 - \theta_0, \theta'_R - \theta_R})^{-1} S_{\delta'_0 - \theta_0, \delta'_R - \theta_R} + (S_{\theta'_0 - \theta_0, \theta'_R - \theta_R})^{-1} (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} \right. \\ &\quad \left. \times S_{\theta_0 - \delta_0, \theta_R - \delta_R} \Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z) S_{\delta'_0 - \delta_0, \delta'_R - \delta_R} \right]^{-1}. \end{aligned} \quad (4.15)$$

Proof. This follows from multiplying (4.6) by $S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ from the right, inserting the terms $S_{\delta'_0 - \delta_0, \delta'_R - \delta_R} (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1}$ to the right of $\Lambda_{\delta_0, \delta_R}^{\delta'_0, \delta'_R}(z)$ twice, and then algebraically manipulate the various terms in the equation resulting from these insertions into (4.6) (repeatedly using $(ST)^{-1} = T^{-1}S^{-1}$, etc.). \square

Since the self-adjoint case is the principal focus for the remainder of this section, we now recall some additional pertinent facts from [52] (see also [42, Sect. 6]) on linear fractional transformations of matrices. Defining

$$J_4 = \begin{bmatrix} 0 & -I_2 \\ I_2 & 0 \end{bmatrix}, \quad (4.16)$$

and

$$\mathcal{A}_4 = \{A \in \mathbb{C}^{4 \times 4} \mid A^* J_4 A = J_4\}, \quad (4.17)$$

representing $A \in \mathbb{C}^{4 \times 4}$ by

$$A = \begin{bmatrix} A_{1,1} & A_{1,2} \\ A_{2,1} & A_{2,2} \end{bmatrix}, \quad A_{p,q} \in \mathbb{C}^{2 \times 2}, \quad 1 \leq p, q \leq 2, \quad (4.18)$$

the condition $A^* J_4 A = J_4$ in (4.17) is equivalent to

$$\begin{aligned} A_{1,1}^* A_{2,1} &= A_{2,1}^* A_{1,1}, & A_{2,2}^* A_{1,2} &= A_{1,2}^* A_{2,2}, \\ A_{2,2}^* A_{1,1} - A_{1,2}^* A_{2,1} &= I_2 = A_{1,1}^* A_{2,2} - A_{2,1}^* A_{1,2}, \end{aligned} \quad (4.19)$$

or equivalently, to

$$\begin{bmatrix} A_{2,2}^* & -A_{1,2}^* \\ -A_{2,1}^* & A_{1,1}^* \end{bmatrix} \begin{bmatrix} A_{1,1} & A_{1,2} \\ A_{2,1} & A_{2,2} \end{bmatrix} = I_4. \quad (4.20)$$

Since left inverses in $\mathbb{C}^{4 \times 4}$ are also right inverses, (4.20) implies

$$\begin{bmatrix} A_{1,1} & A_{1,2} \\ A_{2,1} & A_{2,2} \end{bmatrix} \begin{bmatrix} A_{2,2}^* & -A_{1,2}^* \\ -A_{2,1}^* & A_{1,1}^* \end{bmatrix} = I_4, \quad (4.21)$$

that is,

$$\begin{aligned} A_{1,1}A_{1,2}^* &= A_{1,2}A_{1,1}^*, & A_{2,2}A_{2,1}^* &= A_{2,1}A_{2,2}^*, \\ A_{2,2}A_{1,1}^* - A_{2,1}A_{1,2}^* &= I_2 = A_{1,1}A_{2,2}^* - A_{1,2}A_{2,1}^*, \end{aligned} \quad (4.22)$$

or equivalently,

$$AJ_4A^* = J_4. \quad (4.23)$$

In particular,

$$A \in \mathcal{A}_4 \text{ if and only if } A^{-1} \in \mathcal{A}_4. \quad (4.24)$$

At this point we turn to the particularly important special self-adjoint case where V and $\theta_0, \theta_R, \theta'_0, \theta'_R$ are real-valued. In this case we will now prove that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ is a 2×2 matrix-valued Herglotz function. But first we note the following result:

Lemma 4.4. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R, \delta_0, \delta_R, \delta'_0, \delta'_R \in [0, 2\pi)$, $\theta'_0 - \theta_0 \neq 0 \pmod{\pi}$, $\theta'_R - \theta_R \neq 0 \pmod{\pi}$, $\delta'_0 - \delta_0 \neq 0 \pmod{\pi}$, $\delta'_R - \delta_R \neq 0 \pmod{\pi}$, and introduce in accordance with (4.15),*

$$\begin{aligned} A(\theta, \delta) &= [A(\theta, \delta)_{j,k}]_{1 \leq j, k \leq 2} \in \mathbb{C}^{4 \times 4}, \\ A(\theta, \delta)_{1,1} &= (S_{\theta'_0 - \theta_0, \theta'_R - \theta_R})^{-1} S_{\delta'_0 - \delta_0, \delta'_R - \delta_R}, \\ A(\theta, \delta)_{1,2} &= (S_{\theta'_0 - \theta_0, \theta'_R - \theta_R})^{-1} (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\theta_0 - \delta_0, \theta_R - \delta_R}, \\ A(\theta, \delta)_{2,1} &= S_{\delta'_0 - \theta'_0, \delta'_R - \theta'_R}, \\ A(\theta, \delta)_{2,2} &= (S_{\delta'_0 - \delta_0, \delta'_R - \delta_R})^{-1} S_{\theta'_0 - \delta_0, \theta'_R - \delta_R}. \end{aligned} \quad (4.25)$$

Then

$$A(\theta, \delta) \in \mathcal{A}_4. \quad (4.26)$$

Proof. Since according to (2.43) $S_{\alpha, \beta}$ are 2×2 diagonal matrices, an entirely straightforward (though, admittedly, rather tedious) explicit computation shows that the block matrix entries of $A_{\theta, \delta}$ in (4.25) satisfy the relations in (4.19). \square

We denote by \mathbb{C}_+ the open complex upper half-plane and abbreviate $\text{Im}(L) = (L - L^*)/(2i)$ for $L \in \mathbb{C}^{n \times n}$, $n \in \mathbb{N}$. In addition, $d\|\Sigma\|_{\mathbb{C}^2}$ will denote the total variation of the 2×2 matrix-valued measure $d\Sigma$ below in (4.32).

We recall that $M(\cdot)$ is called an $n \times n$ matrix-valued Herglotz function if it is analytic on \mathbb{C}_+ and $\text{Im}(M(z)) \geq 0$ for all $z \in \mathbb{C}_+$. In this context we also recall the following result:

Lemma 4.5. *Assume that $A = [A_{j,k}]_{1 \leq j,k \leq 2} \in \mathcal{A}_4$ and $L \in \mathbb{C}^{2 \times 2}$. Then*

$$\operatorname{Im}(L) > 0 \text{ implies } \ker(A_{1,1} + A_{1,2}L) = \{0\} \quad (4.27)$$

and $M_A(L) = (A_{2,1} + A_{2,2}L)(A_{1,1} + A_{1,2}L)^{-1}$ (defined according to (4.1)) satisfies

$$\operatorname{Im}(M_A) = ((A_{2,1} + A_{2,2}L)^{-1})^* \operatorname{Im}(L)(A_{2,1} + A_{2,2}L)^{-1} > 0. \quad (4.28)$$

In particular, if $M(\cdot)$ is a 2×2 matrix-valued Herglotz function satisfying

$$\operatorname{Im}(M(z)) > 0, \quad z \in \mathbb{C}_+, \quad (4.29)$$

then $M_A(\cdot)$ (in obvious notation defined according to (4.1) with L replaced by $M(\cdot)$) is a 2×2 matrix-valued Herglotz function satisfying

$$\operatorname{Im}(M_A(z)) > 0, \quad z \in \mathbb{C}_+. \quad (4.30)$$

Proof. This is the special finite-dimensional case of [42, Theorem 6.4]. \square

Now we are in position, in the case where H_{θ_0, θ_R} is self-adjoint, to prove the fundamental Herglotz property of the matrix $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$.

Theorem 4.6. *Let $\theta_0, \theta_R, \theta'_0, \theta'_R \in [0, 2\pi)$, $\theta'_0 - \theta_0 \neq 0 \pmod{\pi}$, $\theta'_R - \theta_R \neq 0 \pmod{\pi}$, $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$, and H_{θ_0, θ_R} be defined as in (2.5). In addition, suppose that V is real-valued (and hence H_{θ_0, θ_R} is self-adjoint). Then $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ is a 2×2 matrix-valued Herglotz function admitting the representation*

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} = \Xi_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} + \int_{\mathbb{R}} d\Sigma_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\lambda) \left(\frac{1}{\lambda - z} - \frac{\lambda}{1 + \lambda^2} \right), \quad (4.31)$$

$$z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}),$$

$$\Xi_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} = \left(\Xi_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} \right)^* \in \mathbb{C}^{2 \times 2}, \quad \int_{\mathbb{R}} \frac{d \|\Sigma_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\lambda)\|_{\mathbb{C}^2}}{1 + \lambda^2} < \infty, \quad (4.32)$$

where

$$\Sigma_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}((\lambda_1, \lambda_2]) = \frac{1}{\pi} \lim_{\delta \downarrow 0} \lim_{\varepsilon \downarrow 0} \int_{\lambda_1 + \delta}^{\lambda_2 + \delta} d\lambda \operatorname{Im} \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\lambda + i\varepsilon) \right), \quad (4.33)$$

$$\lambda_1, \lambda_2 \in \mathbb{R}, \quad \lambda_1 < \lambda_2.$$

In addition,

$$\operatorname{Im} \left(\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} \right) > 0, \quad z \in \mathbb{C}_+. \quad (4.34)$$

Proof. Without loss of generality, we take $z \in \mathbb{C}_+$. An analytic continuation with respect to z then extends the result (4.31) to $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$.

We will first prove the Herglotz property for $\Lambda_{\pi/2, \pi/2}$ and then use a special case of the linear fractional transformation (4.15) to conclude that $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$ is a 2×2 matrix-valued Herglotz function for any $\theta_0, \theta_R, \theta'_0, \theta'_R \in [0, 2\pi)$. Our point of departure is formula (3.40) in the special case $\theta_0 = \theta_R = \pi/2$, that is,

$$\Lambda_{\pi/2, \pi/2}(z) = \gamma_{0,0} [\gamma_{0,0} (H_{\pi/2, \pi/2} - \bar{z})^{-1}]^*, \quad z \in \mathbb{C}_+, \quad (4.35)$$

noticing that

$$\widehat{\gamma}_{\pi/2, \pi/2} = \gamma_{\pi, \pi} = -\gamma_{0,0} \quad (4.36)$$

by (2.6).

First, we slightly change the definition of $\gamma_{0,0}$ by introducing

$$\widetilde{\gamma}_{0,0}: \begin{cases} H^1((0, R)) \rightarrow \mathbb{C}^2, \\ u \mapsto \begin{bmatrix} u(0) \\ u(R) \end{bmatrix}, \end{cases} \quad \widetilde{\gamma}_{0,0} \in \mathcal{B}(H^1((0, R)), \mathbb{C}^2), \quad (4.37)$$

instead. One notes that $\widetilde{\gamma}_{0,0}$ is well-defined in the following sense: Any $u \in H^1((0, R))$ has a representative in its equivalence class of Lebesgue measurable and square integrable elements, again denoted by u for simplicity, that is absolutely continuous on $[0, R]$. In fact, by a standard Sobolev embedding result, one has $H^1(0, R) \hookrightarrow C^{1/2}((0, R)) = C^{1/2}([0, R])$. In particular, the limits $\lim_{x \downarrow 0} u(x) = u(0)$ and $\lim_{x \uparrow R} u(x) = u(R)$ are well-defined for this representative u .

Next, using

$$H^1((0, R)) \hookrightarrow L^2((0, R); dx) \hookrightarrow H^1((0, R))^*, \quad (4.38)$$

one infers that (cf. also Remark 3.7)

$$(\widetilde{\gamma}_{0,0})^*: \begin{cases} \mathbb{C}^2 \rightarrow H^1((0, R))^*, \\ \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \mapsto v_1 \delta_0 + v_2 \delta_R, \end{cases} \quad (\widetilde{\gamma}_{0,0})^* \in \mathcal{B}(\mathbb{C}^2, H^1((0, R))^*), \quad (4.39)$$

where (cf. also [48, Example 2])

$$\begin{aligned} \delta_0, \delta_R &\in H^1((0, R))^*, \\ \delta_0(u) &= {}_{H^1((0, R))} \langle \bar{u}, \delta_0 \rangle_{H^1((0, R))^*} = u(0), \\ \delta_R(u) &= {}_{H^1((0, R))} \langle \bar{u}, \delta_R \rangle_{H^1((0, R))^*} = u(R), \quad u \in H^1((0, R)), \end{aligned} \quad (4.40)$$

with ${}_{H^1((0, R))} \langle \cdot, \cdot \rangle_{H^1((0, R))^*}$ denoting the duality pairing between $H^1((0, R))$ and $H^1((0, R))^*$ (linear in the second argument and antilinear in the first, see, e.g., [37, Sect. 2] for more details). Indeed, (4.40) follows from

$$u(0)v_1 + u(R)v_2 = (\widetilde{\gamma}_{0,0} \bar{u}, v)_{\mathbb{C}^2} = {}_{H^1((0, R))} \langle \bar{u}, (\widetilde{\gamma}_{0,0})^* v \rangle_{H^1((0, R))^*}, \quad (4.41)$$

$$v = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \in \mathbb{C}^2, \quad u \in H^1((0, R)). \quad (4.42)$$

Next, still following the material discussed in [37, Sect. 2], we extend the operator $H_{\pi/2, \pi/2}$ in $L^2((0, R); dx)$,

$$(H_{\pi/2, \pi/2} - zI) : \text{dom}(H_{\pi/2, \pi/2}) \rightarrow L^2((0, R); dx), \quad z \in \mathbb{C}_+, \quad (4.43)$$

where $\text{dom}(H_{\pi/2, \pi/2}) \hookrightarrow L^2((0, R); dx)$, to its extension $\widetilde{H_{\pi/2, \pi/2}}$, which maps $H^1((0, R))$ boundedly into $H^1((0, R))^*$,

$$\widetilde{H_{\pi/2, \pi/2}} \in \mathcal{B}(H^1((0, R)), H^1((0, R))^*), \quad (4.44)$$

such that (with $\widetilde{I} : H^1((0, R)) \hookrightarrow H^1((0, R))^*$ the continuous embedding operator)

$$(\widetilde{H_{\pi/2, \pi/2}} + \widetilde{I}) \in \mathcal{B}(H^1((0, R)), H^1((0, R))^*) \quad (4.45)$$

$$\text{and } (\widetilde{H_{\pi/2, \pi/2}} + \widetilde{I}) : H^1((0, R)) \rightarrow H^1((0, R))^* \text{ is unitary.} \quad (4.46)$$

In addition (cf. (2.19)),

$$(H_{\pi/2, \pi/2} + I)^{1/2} \in \mathcal{B}(H^1((0, R)), L^2((0, R); dx)) \quad (4.47)$$

$$\text{and } (H_{\pi/2, \pi/2} + I)^{1/2} : H^1((0, R)) \rightarrow L^2((0, R); dx) \text{ is unitary} \quad (4.48)$$

(cf. [37, Sect. 2]). Moreover, $\widetilde{H_{\pi/2, \pi/2}}$ is self-adjoint,

$$\left(\widetilde{H_{\pi/2, \pi/2}}\right)^* = \widetilde{H_{\pi/2, \pi/2}}, \quad (4.49)$$

in the sense that

$$H^1((0, R)) \langle w_1, \widetilde{H_{\pi/2, \pi/2}} w_2 \rangle_{H^1((0, R))^*} = \overline{H^1((0, R)) \langle w_2, \widetilde{H_{\pi/2, \pi/2}} w_1 \rangle_{H^1((0, R))^*}}, \quad (4.50)$$

$w_1, w_2 \in H^1((0, R))$

(again we refer to [37, Sect. 2] for more details).

In addition,

$$\left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I}\right)^{-1} : H^1((0, R))^* \rightarrow H^1((0, R)), \quad z \in \mathbb{C}_+, \quad (4.51)$$

and hence,

$$\left(\left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I}\right)^{-1} w\right)(x) = H^1((0, R)) \langle \overline{G_{\pi/2, \pi/2}(z, x, \cdot)}, w \rangle_{H^1((0, R))^*}, \quad (4.52)$$

$w \in H^1((0, R))^*, z \in \mathbb{C}_+$

using the fact that

$$G_{\pi/2, \pi/2}(z, x, \cdot) \in H^1((0, R)), \quad x \in \mathbb{R}, z \in \mathbb{C} \setminus \sigma(H_{\pi/2, \pi/2}). \quad (4.53)$$

By (2.34) and (2.35), the Wronskian W is of the form

$$\begin{aligned} W(z) &= W(u_{+,\pi/2}(z, \cdot), u_{-,\pi/2}(z, \cdot)) \\ &= -u'_{+,\pi/2}(z, 0)u_{-,\pi/2}(z, 0) = u_{+,\pi/2}(z, R)u'_{-,\pi/2}(z, R). \end{aligned} \quad (4.54)$$

Using (3.37), one computes

$$\begin{aligned} &\left(\left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} (\widetilde{\gamma}_{0,0})^* v \right)(x) \\ &= \frac{1}{W(z)} {}_{H^1((0,R))} \langle \overline{G_{\pi/2, \pi/2}(z, x, \cdot)}, [v_1 \delta_0 + v_2 \delta_R] \rangle_{{}_{H^1((0,R))}^*} \\ &= \frac{1}{W(z)} (u_{-,\pi/2}(z, x)u_{+,\pi/2}(z, R)v_2 + u_{+,\pi/2}(z, x)u_{-,\pi/2}(z, 0)v_1), \end{aligned} \quad (4.55)$$

$$v = [v_1 \ v_2]^\top \in \mathbb{C}^2,$$

and hence

$$\begin{aligned} &\widetilde{\gamma}_{0,0} \left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} (\widetilde{\gamma}_{0,0})^* v \\ &= \frac{1}{W(z)} \begin{bmatrix} u_{-,\pi/2}(z, 0)u_{+,\pi/2}(z, R)v_2 + u_{-,\pi/2}(z, 0)v_1 \\ u_{+,\pi/2}(z, R)v_2 + u_{+,\pi/2}(z, R)u_{-,\pi/2}(z, 0)v_1 \end{bmatrix} \\ &= \frac{1}{W(z)} \begin{bmatrix} u_{-,\pi/2}(z, 0) & u_{-,\pi/2}(z, 0)u_{+,\pi/2}(z, R) \\ u_{-,\pi/2}(z, 0)u_{+,\pi/2}(z, R) & u_{+,\pi/2}(z, R) \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \\ &= \Lambda_{\pi/2, \pi/2}(z)v, \quad v = [v_1 \ v_2]^\top \in \mathbb{C}^2, \end{aligned} \quad (4.56)$$

inserting (4.54) for $W(\cdot)$. Consequently, one has

$$\Lambda_{\pi/2, \pi/2}(z) = \widetilde{\gamma}_{0,0} \left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} (\widetilde{\gamma}_{0,0})^*, \quad z \in \mathbb{C}_+, \quad (4.57)$$

and also

$$\Lambda_{\pi/2, \pi/2}(z)^* = \Lambda_{\pi/2, \pi/2}(\bar{z}), \quad z \in \mathbb{C}_+. \quad (4.58)$$

In particular,

$$\begin{aligned} (v, \operatorname{Im}(\Lambda_{\pi/2, \pi/2}(z))v)_{\mathbb{C}^2} &= \frac{1}{2i} (v, (\Lambda_{\pi/2, \pi/2}(z) - \Lambda_{\pi/2, \pi/2}(z)^*)v)_{\mathbb{C}^2} \\ &= \frac{1}{2i} (v, \widetilde{\gamma}_{0,0} \left(\left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} - \left(\widetilde{H_{\pi/2, \pi/2}} - \bar{z}\widetilde{I} \right)^{-1} \right) (\widetilde{\gamma}_{0,0})^* v)_{\mathbb{C}^2} \\ &= \operatorname{Im}(z) (v, \widetilde{\gamma}_{0,0} \left[\left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} \widetilde{I} \left(\widetilde{H_{\pi/2, \pi/2}} - \bar{z}\widetilde{I} \right)^{-1} \right] (\widetilde{\gamma}_{0,0})^* v)_{\mathbb{C}^2} \\ &= \operatorname{Im}(z) (v, \left[\widetilde{\gamma}_{0,0} \left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} \right] \widetilde{I} \left[\widetilde{\gamma}_{0,0} \left(\widetilde{H_{\pi/2, \pi/2}} - z\widetilde{I} \right)^{-1} \right]^* v)_{\mathbb{C}^2} \end{aligned}$$

$$\begin{aligned}
&= \operatorname{Im}(z) {}_{H^1((0,R))} \left\langle \left[\widetilde{\gamma}_{0,0} \left(\widetilde{H}_{\pi/2,\pi/2} - z\widetilde{I} \right)^{-1} \right]^* v, \widetilde{I} \right. \\
&\quad \left. \times \left[\widetilde{\gamma}_{0,0} \left(\widetilde{H}_{\pi/2,\pi/2} - z\widetilde{I} \right)^{-1} \right]^* v \right\rangle_{H^1((0,R))^*} \\
&= \operatorname{Im}(z) \left(\left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* v, \left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* v \right)_{L^2((0,R);dx)} \\
&= \operatorname{Im}(z) \left\| \left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* v \right\|_{L^2((0,R);dx)}^2 \geq 0, \quad v \in \mathbb{C}^2, z \in \mathbb{C}_+, \tag{4.59}
\end{aligned}$$

since the duality pairing ${}_{H^1((0,R))} \langle \cdot, \cdot \rangle_{H^1((0,R))^*}$ between the spaces $H^1((0,R))$ and $H^1((0,R))^*$ is compatible with the scalar product in $L^2((0,R);dx)$, that is,

$${}_{H^1((0,R))} \langle w_1, \widetilde{I} w_2 \rangle_{H^1((0,R))^*} = (w_1, w_2)_{L^2((0,R);dx)}, \quad w_1, w_2 \in H^1((0,R)). \tag{4.60}$$

In particular,

$$\operatorname{Im}(\Lambda_{\pi/2,\pi/2}(z)) = \left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right] \left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* \geq 0, \tag{4.61}$$

$$z \in \mathbb{C}_+,$$

and thus, $\Lambda_{\pi/2,\pi/2}(\cdot)$ is a 2×2 matrix-valued Herglotz function. In fact, one can improve on (4.61) to obtain

$$\operatorname{Im}(\Lambda_{\pi/2,\pi/2}(z)) > 0, \quad z \in \mathbb{C}_+, \tag{4.62}$$

since

$$\ker \left(\left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* \right) = \{0\}, \quad z \in \mathbb{C} \setminus \sigma(H_{\pi/2,\pi/2}). \tag{4.63}$$

To prove (4.63), one can argue as follows: Suppose that

$$[a_0 \ a_R]^\top \in \ker \left(\left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* \right), \tag{4.64}$$

then by (3.54),

$$\begin{aligned}
&\left(\left[\gamma_{0,0} (H_{\pi/2,\pi/2} - zI)^{-1} \right]^* [a_0 \ a_R]^\top \right)(x) = \frac{1}{W(u_{+,\pi/2}(\bar{z}, \cdot), u_{-,\pi/2}(\bar{z}, \cdot))} \\
&\quad \times [a_0 u_{-,\pi/2}(\bar{z}, 0) u_{+,\pi/2}(\bar{z}, x) + a_R u_{+,\pi/2}(\bar{z}, R) u_{-,\pi/2}(\bar{z}, x)] = 0, \tag{4.65} \\
&\quad z \in \mathbb{C} \setminus \sigma(H_{\pi/2,\pi/2}).
\end{aligned}$$

Since by definition (cf. (2.34), (2.35)), $u'_{-,\pi/2}(\bar{z}, 0) = u'_{+,\pi/2}(\bar{z}, R) = 0$, one concludes that

$$u_{-,\pi/2}(\bar{z}, 0) \neq 0, \quad u_{+,\pi/2}(\bar{z}, R) \neq 0. \tag{4.66}$$

Moreover, since $W(u_{+,\pi/2}(\bar{z}, \cdot), u_{-,\pi/2}(\bar{z}, \cdot)) \neq 0$ for all $z \in \mathbb{C} \setminus \sigma(H_{\pi/2,\pi/2})$ (otherwise, \bar{z} would be an eigenvalue of $H_{\pi/2,\pi/2}$), $u_{+,\pi/2}(\bar{z}, \cdot)$ and $u_{-,\pi/2}(\bar{z}, \cdot)$ are linearly independent, implying

$$a_0 = a_R = 0, \tag{4.67}$$

and hence (4.63).

Next, using the notation introduced in (3.1), and applying the linear fractional transformation (4.6) one can show that (with $z \in \mathbb{C}_+$)

$$\Lambda_{\pi/2, \pi/2}(z) [C_{\pi/2, \pi/2} + S_{\pi/2, \pi/2} \Lambda_{0,0}(z)] = [C_{\pi, \pi} + S_{\pi, \pi} \Lambda_{0,0}(z)], \quad (4.68)$$

or equivalently, that

$$\Lambda_{\pi/2, \pi/2}(z) \Lambda_{0,0}(z) = -I_2, \quad (4.69)$$

and hence,

$$\Lambda_{0,0}^{\pi/2, \pi/2}(z) = \Lambda_{0,0}(z) = -[\Lambda_{\pi/2, \pi/2}(z)]^{-1} = -[\Lambda_{\pi/2, \pi/2}^{\pi, \pi}(z)]^{-1}, \quad (4.70)$$

is a 2×2 matrix-valued Herglotz function too, satisfying

$$\operatorname{Im}(\Lambda_{0,0}^{\pi/2, \pi/2}(z)) > 0, \quad z \in \mathbb{C}_+. \quad (4.71)$$

The Herglotz property of $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$, $\theta_0, \theta_R, \theta'_0, \theta'_R \in [0, 2\pi)$ then follows again from the linear fractional transformation (4.15) and from (4.30) in Lemma 4.5 upon identifying the 2×2 block matrix $A = [A_{j,k}]_{1 \leq j, k \leq 2}$ in Lemma 4.5 with $A(\theta, \delta) \in \mathcal{A}_4$ in the special case where

$$\begin{aligned} \theta_0, \theta_R, \theta'_0, \theta'_R &\in [0, 2\pi), \quad \theta'_0 - \theta_0 \neq 0 \pmod{\pi}, \quad \theta'_R - \theta_R \neq 0 \pmod{\pi}, \\ \delta_0 = \delta_R &= 0, \quad \delta'_0 = \delta'_R = \pi/2. \end{aligned} \quad (4.72)$$

Given the Herglotz property of $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}$, one obtains as in [42, Theorem 5.4] the representation

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} = \Xi_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} + \Upsilon_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} z + \int_{\mathbb{R}} d\Sigma_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\lambda) \left(\frac{1}{\lambda - z} - \frac{\lambda}{1 + \lambda^2} \right), \quad (4.73)$$

$$z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}),$$

with $\Xi_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ and $\Sigma_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(\cdot)$ as in (4.32) and (4.33), and with $\Upsilon_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}$ satisfying

$$0 \leq \Upsilon_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} \in \mathbb{C}^{2 \times 2}. \quad (4.74)$$

Thus, to conclude the proof of (4.31), it remains to prove that actually, $\Upsilon_{\theta_0, \theta_R}^{\theta'_0, \theta'_R} = 0$. The latter fact is clear since

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \underset{\substack{|z| \rightarrow \infty \\ \operatorname{Im}(z^{1/2}) > 0}}{=} O(|z|^{1/2}), \quad (4.75)$$

using the fact that by (4.6),

$$\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) = C_{\theta'_0 - \theta_0, \theta'_R - \theta_R} + S_{\theta'_0 - \theta_0, \theta'_R - \theta_R} \Lambda_{\theta_0, \theta_R}(z), \quad (4.76)$$

and applying Lemma 3.4.

Finally, (4.34) is a consequence of (4.71) and Lemma 4.5 with $A = A(\theta, \delta)$ chosen as in (4.72). \square

As a particular case of Theorem 4.6, for the generalized Dirichlet-to-Neumann map, where $\theta'_0 = (\theta_0 + (\pi/2)) \bmod(2\pi)$, $\theta'_R = (\theta_R + (\pi/2)) \bmod(2\pi)$, one concludes that

$$\Lambda_{\theta_0, \theta_R}(z) = \Lambda_{\theta_0, \theta_R}^{(\theta_0 + \frac{\pi}{2}) \bmod(2\pi), (\theta_R + \frac{\pi}{2}) \bmod(2\pi)}(z), \quad \theta_0, \theta_R \in [0, 2\pi), \quad z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R}), \quad (4.77)$$

is a 2×2 matrix-valued Herglotz function.

5. Krein-Type Resolvent Formulas

Krein-type resolvent formulas have been studied in a great variety of contexts, far too numerous to account for all in this paper. For instance, they are of fundamental importance in connection with the spectral and inverse spectral theory of ordinary and partial differential operators. Abstract versions of Krein-type resolvent formulas (see also the brief discussion at the end of our introduction), connected to boundary value spaces (boundary triples) and self-adjoint extensions of closed symmetric operators with equal (possibly infinite) deficiency spaces, have received enormous attention in the literature. In particular, we note that Robin-to-Dirichlet maps in the context of ordinary differential operators reduce to the celebrated (possibly, matrix-valued) Weyl–Titchmarsh function, the basic object of spectral analysis in this context. Since it is impossible to cover the literature in this paper, we refer the reader to the rather extensive recent bibliography in [35], [36], and [38]. Here we just mention, for instance, [1, Sect. 84], [2], [3], [4], [5], [6], [11], [12], [14], [15], [16], [17], [18], [19], [25], [26], [28], [33], [34], [39], [42], [44, Ch. 3], [45], [46, Ch. 13], [49], [50], [51], [53], [54], [55], [58], [61], [63], [64], [65], [66], [67], [69], [70], [71], [72], [73], [75], [74], [77], [78], [79], [82], and the references cited therein.

We start by explicitly computing operators of the type $\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}$ which play a role at various places in this manuscript (cf. Theorem 3.5, Lemma 5.2, and Theorem 5.3):

Lemma 5.1. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, let H_{θ_0, θ_R} be defined as in (2.5), and suppose that $z \in \mathbb{C} \setminus \sigma(H_{\theta_0, \theta_R})$. Then, assuming $f \in L^2((0, R); dx)$, and writing*

$$\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}f = \begin{bmatrix} (\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1})_1 f \\ (\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1})_2 f \end{bmatrix} \in \mathbb{C}^2, \quad (5.1)$$

one has

$$\begin{aligned} (\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1})_1 f &= \frac{\sin(\theta'_0 - \theta_0)}{W(u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot))} (\overline{u_{+, \theta_R}(z, \cdot)}, f)_{L^2((0, R); dx)} \\ &\times \begin{cases} -\frac{u_{-, \theta_0}(z, 0)}{\sin(\theta_0)}, & \theta_0 \in S_{2\pi} \setminus \{0, \pi\}, \\ \frac{u'_{-, \theta_0}(z, 0)}{\cos(\theta_0)}, & \theta_0 \in S_{2\pi} \setminus \{\pi/2, 3\pi/2\}, \end{cases} \end{aligned} \quad (5.2)$$

$$\begin{aligned}
(\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1})_2 f &= \frac{-\sin(\theta'_R - \theta_R)}{W(u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot))} (\overline{u_{-, \theta_0}(z, \cdot)}, f)_{L^2((0, R); dx)} \\
&\times \begin{cases} \frac{u_{+, \theta_R}(z, R)}{\sin(\theta_R)}, & \theta_R \in S_{2\pi} \setminus \{0, \pi\}, \\ \frac{u'_{+, \theta_R}(z, R)}{\cos(\theta_R)}, & \theta_R \in S_{2\pi} \setminus \{\pi/2, 3\pi/2\}, \end{cases} \quad (5.3)
\end{aligned}$$

in particular,

$$|(\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1})_1 f|_{\theta'_0 \rightarrow \theta_0} = O(\theta'_0 - \theta_0) C_1(z) \|f\|_{L^2((0, R); dx)}, \quad (5.4)$$

$$|(\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1})_2 f|_{\theta'_R \rightarrow \theta_R} = O(\theta'_R - \theta_R) C_2(z) \|f\|_{L^2((0, R); dx)}, \quad (5.5)$$

for some constants $C_j(z) > 0$, $j = 1, 2$, and hence

$$\gamma_{\theta_0, \theta_R}(H_{\theta_0, \theta_R} - zI)^{-1} = 0 \text{ in } \mathcal{B}(L^2((0, R); dx), \mathbb{C}^2), \quad (5.6)$$

$$(\gamma_{\theta_0, \theta_R}(H_{\theta_0, \theta_R} - zI)^{-1})_k = 0 \text{ in } \mathcal{B}(L^2((0, R); dx), \mathbb{C}), \quad k = 1, 2. \quad (5.7)$$

Proof. Employing (3.37) and (3.38) one obtains

$$\begin{aligned}
\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1} f &= \gamma_{\theta'_0, \theta'_R} \left(\int_0^R dx' G_{\theta_0, \theta_R}(z, \cdot, x') f(x') \right) \\
&= \left[\left(\gamma_{\theta'_0, \theta'_R} \left(\int_0^R dx' G_{\theta_0, \theta_R}(z, \cdot, x') f(x') \right) \right)_1 \right] \\
&\quad \left[\left(\gamma_{\theta'_0, \theta'_R} \left(\int_0^R dx' G_{\theta_0, \theta_R}(z, \cdot, x') f(x') \right) \right)_2 \right] \quad (5.8)
\end{aligned}$$

and hence (with $W(z) = W(u_{+, \theta_R}(z, \cdot), u_{-, \theta_0}(z, \cdot))$)

$$\begin{aligned}
&\left(\gamma_{\theta'_0, \theta'_R} \left(\int_0^R dx' G_{\theta_0, \theta_R}(z, \cdot, x') f(x') \right) \right)_1 \\
&= \frac{1}{W(z)} \left\{ \cos(\theta'_0) \left(\int_0^R dx' G_{\theta_0, \theta_R}(z, 0, x') f(x') \right) \right. \\
&\quad \left. + \sin(\theta'_0) \left(\int_0^R dx' \left(\frac{\partial}{\partial x} G_{\theta_0, \theta_R}(z, x, x') \Big|_{x < x'} \right) \Big|_{x=0} f(x') \right) \right\} \\
&= \frac{1}{W(z)} \left\{ [\cos(\theta'_0) u_{-, \theta_0}(z, 0) + \sin(\theta'_0) u'_{-, \theta_0}(z, 0)] \int_0^R dx' u_{+, \theta_R}(z, x') f(x') \right\} \\
&= \frac{1}{W(z)} (\overline{u_{+, \theta_R}(z, \cdot)}, f)_{L^2((0, R); dx)} \\
&\times \begin{cases} [\cos(\theta'_0) - \sin(\theta'_0) \cot(\theta_0)] u_{-, \theta_0}(z, 0), & \theta_0 \in S_{2\pi} \setminus \{0, \pi\}, \\ [-\cos(\theta'_0) \tan(\theta_0) + \sin(\theta'_0)] u'_{-, \theta_0}(z, 0), & \theta_0 \in S_{2\pi} \setminus \{\pi/2, 3\pi/2\}. \end{cases} \quad (5.9)
\end{aligned}$$

Equation (5.9) is easily seen to be equivalent to (5.2). Equation (5.3) is derived analogously. \square

Introducing the orthogonal projections in \mathbb{C}^2 ,

$$P_1 = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}, \quad P_2 = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}, \quad (5.10)$$

one obtains the following result, patterned after [62, Lemma 6] in the context of Schrödinger operators with Dirichlet and Neumann boundary conditions on a cube in \mathbb{R}^n :

Lemma 5.2. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, let H_{θ_0, θ_R} and $H_{\theta'_0, \theta'_R}$ be defined as in (2.5), and suppose that $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\theta'_0, \theta'_R}))$. Then*

$$\begin{aligned} (H_{\theta'_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\ &+ [\gamma_{\overline{\theta'_0}, \overline{\theta'_R}}((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} [\gamma_{\theta_0, \theta_R}(H_{\theta'_0, \theta'_R} - zI)^{-1}], \end{aligned} \quad (5.11)$$

$\theta'_0 \neq \theta_0, \theta'_R \neq \theta_R,$

$$\begin{aligned} (H_{\theta_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\ &+ [\gamma_{\overline{\theta_0}, \overline{\theta'_R}}((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [\sin(\theta'_R - \theta_R)]^{-1} P_2 [\gamma_{\theta_0, \theta_R}(H_{\theta_0, \theta'_R} - zI)^{-1}], \end{aligned} \quad (5.12)$$

$\theta'_R \neq \theta_R,$

$$\begin{aligned} (H_{\theta'_0, \theta_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\ &+ [\gamma_{\overline{\theta'_0}, \overline{\theta_R}}((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [\sin(\theta'_0 - \theta_0)]^{-1} P_1 [\gamma_{\theta_0, \theta_R}(H_{\theta'_0, \theta_R} - zI)^{-1}], \end{aligned} \quad (5.13)$$

$\theta'_0 \neq \theta_0.$

Proof. We first consider the case

$$\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi} \setminus \{0, \pi\}, \quad \theta'_0 \neq \theta_0, \quad \theta'_R \neq \theta_R, \quad (5.14)$$

which illustrates the principal idea of the proof. To get started, we pick $f, g \in L^2((0, R); dx)$ and introduce

$$\begin{aligned} \phi &= ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1} f \in \text{dom}((H_{\theta_0, \theta_R})^*), \\ \psi &= (H_{\theta'_0, \theta'_R} - zI)^{-1} g \in \text{dom}(H_{\theta'_0, \theta'_R}). \end{aligned} \quad (5.15)$$

Then one computes

$$\begin{aligned} &((H_{\theta_0, \theta_R})^* - \bar{z}I)\phi, \psi)_{L^2((0, R); dx)} - (\phi, (H_{\theta'_0, \theta'_R} - zI)\psi)_{L^2((0, R); dx)} \\ &= - \int_0^R dx \overline{\phi''(x)} \psi(x) + \int_0^R dx \overline{\phi(x)} \psi''(x) \\ &= -\overline{\phi'(R)} \psi(R) + \overline{\phi'(0)} \psi(0) + \overline{\phi(R)} \psi'(R) - \overline{\phi(0)} \psi'(0) \\ &= [\cot(\theta'_R) - \cot(\theta_R)] \overline{\phi(R)} \psi(R) + [\cot(\theta'_0) - \cot(\theta_0)] \overline{\phi(0)} \psi(0) \\ &= -\frac{\sin(\theta'_R - \theta_R)}{\sin(\theta'_R) \sin(\theta_R)} \overline{\phi(R)} \psi(R) - \frac{\sin(\theta'_0 - \theta_0)}{\sin(\theta'_0) \sin(\theta_0)} \overline{\phi(0)} \psi(0), \end{aligned} \quad (5.16)$$

using the fact that (5.15) implies

$$\begin{aligned}
\cos(\bar{\theta}_0)\phi(0) + \sin(\bar{\theta}_0)\phi'(0) &= 0, \\
\cos(\bar{\theta}_R)\phi(R) - \sin(\bar{\theta}_R)\phi'(R) &= 0, \\
\cos(\theta'_0)\psi(0) + \sin(\theta'_0)\psi'(0) &= 0, \\
\cos(\theta'_R)\psi(R) - \sin(\theta'_R)\psi'(R) &= 0.
\end{aligned} \tag{5.17}$$

Using (5.17) once again, one also computes

$$(\gamma_{\bar{\theta}'_0, \bar{\theta}'_R} \phi, \gamma_{\theta_0, \theta_R} \psi)_{\mathbb{C}^2} = -\frac{\sin^2(\theta'_R - \theta_R)}{\sin(\theta'_R) \sin(\theta_R)} \overline{\phi(R)} \psi(R) - \frac{\sin^2(\theta'_0 - \theta_0)}{\sin(\theta'_0) \sin(\theta_0)} \overline{\phi(0)} \psi(0). \tag{5.18}$$

A comparison of (5.16) and (5.18) then yields

$$\begin{aligned}
&((H_{\theta_0, \theta_R})^* - \bar{z}I)\phi, \psi)_{L^2((0, R); dx)} - (\phi, (H_{\theta'_0, \theta'_R} - zI)\psi)_{L^2((0, R); dx)} \\
&= (\gamma_{\bar{\theta}'_0, \bar{\theta}'_R} \phi, S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} \gamma_{\theta_0, \theta_R} \psi)_{\mathbb{C}^2},
\end{aligned} \tag{5.19}$$

or equivalently,

$$\begin{aligned}
&(f, (H_{\theta'_0, \theta'_R} - zI)^{-1}g)_{L^2((0, R); dx)} = (f, (H_{\theta_0, \theta_R} - zI)^{-1}g)_{L^2((0, R); dx)} \\
&+ (f, [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} [\gamma_{\theta_0, \theta_R} (H_{\theta'_0, \theta'_R} - zI)^{-1}]g)_{L^2((0, R); dx)},
\end{aligned} \tag{5.20}$$

and hence (5.11) since $f, g \in L^2((0, R); dx)$ are arbitrary, under the additional assumptions in (5.14).

Employing Lemma 5.1 (and particularly, (5.4)–(5.7)) then shows that (5.11) is continuous in $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$ with respect to the norm in $\mathcal{B}(L^2((0, R); dx))$, removing the restrictions $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi} \setminus \{0, \pi\}$ in (5.20), and thus proving (5.11).

Analogous considerations imply (5.12) and (5.13). \square

The principal result of this section, Krein's formula for the difference of resolvents of $H_{\theta'_0, \theta'_R}$ and H_{θ_0, θ_R} , then reads as follows:

Theorem 5.3. *Assume that $\theta_0, \theta_R, \theta'_0, \theta'_R \in S_{2\pi}$, let H_{θ_0, θ_R} and $H_{\theta'_0, \theta'_R}$ be defined as in (2.5), and suppose that $z \in \mathbb{C} \setminus (\sigma(H_{\theta_0, \theta_R}) \cup \sigma(H_{\theta'_0, \theta'_R}))$. Then, with $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$ introduced in (2.39),*

$$\begin{aligned}
(H_{\theta'_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\
&- [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R} ((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} \\
&\times \left[\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right]^{-1} [\gamma_{\theta_0, \theta_R} (H_{\theta_0, \theta_R} - zI)^{-1}], \quad \theta_0 \neq \theta'_0, \theta_R \neq \theta'_R.
\end{aligned} \tag{5.21}$$

$$\begin{aligned}
(H_{\theta_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\
&\quad - [\gamma_{\bar{\theta}_0, \bar{\theta}'_R}((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [\sin(\theta'_R - \theta_R)]^{-1} P_2 \\
&\quad \times \left[\Lambda_{\theta_0, \theta'_R}^{\theta_0, \theta'_R}(z) \right]^{-1} P_2 [\gamma_{\theta_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \quad \theta_R \neq \theta'_R,
\end{aligned} \tag{5.22}$$

$$\begin{aligned}
(H_{\theta'_0, \theta_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\
&\quad - [\gamma_{\bar{\theta}'_0, \bar{\theta}_R}((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* [\sin(\theta'_0 - \theta_0)]^{-1} P_1 \\
&\quad \times \left[\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta_R}(z) \right]^{-1} P_1 [\gamma_{\theta'_0, \theta_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \quad \theta_0 \neq \theta'_0.
\end{aligned} \tag{5.23}$$

Proof. Taking adjoints on both sides of (5.11), and subsequently replacing z by \bar{z} , θ_0, θ_R by $\bar{\theta}_0, \bar{\theta}_R$, and V by \bar{V} , then yields

$$\begin{aligned}
(H_{\theta'_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\
&\quad + [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R}((H_{\theta'_0, \theta'_R})^* - \bar{z}I)^{-1}]^* S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} [\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \\
&\quad \theta'_0 \neq \theta_0, \theta'_R \neq \theta_R.
\end{aligned} \tag{5.24}$$

Applying $\gamma_{\theta_0, \theta_R}$ to both sides of (5.24), and using the fact that $\gamma_{\theta_0, \theta_R}(H_{\theta_0, \theta_R} - zI)^{-1} = 0$ by (5.6), one obtains

$$\begin{aligned}
\gamma_{\theta_0, \theta_R}(H_{\theta'_0, \theta'_R} - zI)^{-1} &= \gamma_{\theta_0, \theta_R} [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R}((H_{\theta'_0, \theta'_R})^* - \bar{z}I)^{-1}]^* \\
&\quad \times S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} [\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \quad \theta'_0 \neq \theta_0, \theta'_R \neq \theta_R.
\end{aligned} \tag{5.25}$$

An insertion of (5.25) into (5.11) implies

$$\begin{aligned}
(H_{\theta'_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\
&\quad + [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R}((H_{\theta_0, \theta_R})^* - \bar{z}I)^{-1}]^* S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} \gamma_{\theta_0, \theta_R} [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R}((H_{\theta'_0, \theta'_R})^* - \bar{z}I)^{-1}]^* \\
&\quad \times S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}^{-1} [\gamma_{\theta'_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \quad \theta'_0 \neq \theta_0, \theta'_R \neq \theta_R.
\end{aligned} \tag{5.26}$$

Using (2.48) and (3.39) one obtains

$$\gamma_{\theta_0, \theta_R} [\gamma_{\bar{\theta}'_0, \bar{\theta}'_R}((H_{\theta'_0, \theta'_R})^* - \bar{z}I)^{-1}]^* = - \left[\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z) \right]^{-1} S_{\theta'_0 - \theta_0, \theta'_R - \theta_R}, \quad z \in \mathbb{C}_+, \tag{5.27}$$

and inserting (5.27) into (5.26) yields (5.21).

Since equations (5.22) and (5.23) are proved similarly, we just briefly sketch the proof of (5.22): First, in analogy to (5.24), one derives from (5.12) that

$$\begin{aligned}
(H_{\theta_0, \theta'_R} - zI)^{-1} &= (H_{\theta_0, \theta_R} - zI)^{-1} \\
&\quad + [\gamma_{\bar{\theta}_0, \bar{\theta}'_R}((H_{\theta_0, \theta'_R})^* - \bar{z}I)^{-1}]^* [\sin(\theta'_R - \theta_R)]^{-1} P_2 [\gamma_{\theta_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \\
&\quad \theta'_R \neq \theta_R.
\end{aligned} \tag{5.28}$$

Applying $\gamma_{\theta_0, \theta_R}$ to both sides of (5.28), and using again the fact that $\gamma_{\theta_0, \theta_R}(H_{\theta_0, \theta_R} - zI)^{-1} = 0$, one obtains

$$\begin{aligned} \gamma_{\theta_0, \theta_R}(H_{\theta_0, \theta'_R} - zI)^{-1} &= \gamma_{\theta_0, \theta_R} [\gamma_{\overline{\theta_0, \theta_R}}((H_{\theta_0, \theta'_R})^* - \bar{z}I)^{-1}]^* \\ &\times [\sin(\theta'_R - \theta_R)]^{-1} P_2 [\gamma_{\theta_0, \theta'_R}(H_{\theta_0, \theta_R} - zI)^{-1}], \quad \theta'_R \neq \theta_R. \end{aligned} \quad (5.29)$$

Finally, an insertion of (5.29) into the right-hand side of (5.12), and using (5.27) in the special case $\theta'_0 = \theta_0$, that is,

$$\begin{aligned} \gamma_{\theta_0, \theta_R} [\gamma_{\overline{\theta_0, \theta_R}}((H_{\theta_0, \theta'_R})^* - \bar{z}I)^{-1}]^* &= -[\Lambda_{\theta_0, \theta'_R}^{\theta_0, \theta'_R}(z)]^{-1} S_{0, \theta'_R - \theta_R} \\ &= -[\Lambda_{\theta_0, \theta'_R}^{\theta_0, \theta'_R}(z)]^{-1} [\sin(\theta'_R - \theta_R)]^{-1} P_2, \quad z \in \mathbb{C}_+, \end{aligned} \quad (5.30)$$

yields (5.22). □

Connections between boundary data maps, perturbation determinants, and trace formulas for the resolvent differences described in Theorem 5.3 are worked out in [43].

6. A Brief Outlook

In this section we provide a brief comparison between $\Lambda_{\theta_0, \theta_R}(z)$ and the 2×2 matrix-valued Weyl–Titchmarsh function associated with H_{θ_0, θ_R} in the self-adjoint context, that is, under the assumptions $\theta_0, \theta_R \in [0, 2\pi)$ and V is real-valued in addition to (2.4). While both 2×2 matrices are matrix-valued Herglotz functions, they are quite different as this brief section will show. A more detailed discussion of the interrelations between these two matrices is beyond the scope of this paper and will be taken up elsewhere.

To illustrate some differences between $\Lambda_{\theta_0, \theta_R}(z)$ and various versions of the Weyl–Titchmarsh matrix we will link the entries in both matrices to the Green’s function $G_{\theta_0, \theta_R}(z, x, x')$ of H_{θ_0, θ_R} and provide some formulas which may well be of independent interest.

We start by linking $\Lambda_{\theta_0, \theta_R}(z)$ and $G_{\theta_0, \theta_R}(z, x, x')$ and list a variety of pertinent formulas (choos-

ing $z \in \mathbb{C} \setminus \mathbb{R}$ for notational simplicity):

$$\begin{aligned}
\Lambda_{\theta_0, \theta_R}(z)_{1,1} &= m_{+, \theta_0}(z, \theta_R) \\
&= \frac{-\sin(\theta_0) + \cos(\theta_0)m_{+,0}(z, \theta_R)}{\cos(\theta_0) + \sin(\theta_0)m_{+,0}(z, \theta_R)} \\
&= \frac{-\sin(\theta_0) + \cos(\theta_0)u'_{+, \theta_R}(z, 0)}{\cos(\theta_0) + \sin(\theta_0)u'_{+, \theta_R}(z, 0)} \\
&= \frac{1}{\sin^2(\theta_0)} [G_{\theta_0, \theta_R}(z, 0, 0) + \sin(\theta_0) \cos(\theta_0)], \tag{6.1}
\end{aligned}$$

$$\theta_0 \in [0, 2\pi) \setminus \{0, \pi\}, \theta_R \in [0, 2\pi),$$

$$\begin{aligned}
\Lambda_{0,0}(z)_{1,1} &= m_{+,0}(z, 0) \\
&= \lim_{0 < x < x', x' \downarrow 0} \partial_x \partial_{x'} G_{0,0}(z, x, x'). \tag{6.2}
\end{aligned}$$

Here we used (3.25) and (cf. also (3.24))

$$\begin{aligned}
G_{\theta_0, \theta_R}(z, 0, 0) &= \frac{-\sin(\theta_0)}{\cos(\theta_0) + \sin(\theta_0)m_{+,0}(z, \theta_R)} \\
&= \sin(\theta_0) [-\cos(\theta_0) + \sin(\theta_0)m_{+, \theta_0}(z, \theta_R)]. \tag{6.3}
\end{aligned}$$

In the same manner one obtains

$$\begin{aligned}
\Lambda_{\theta_0, \theta_R}(z)_{2,2} &= -m_{-, \theta_R}(z, \theta_0) \\
&= -\frac{\sin(\theta_R) + \cos(\theta_R)m_{-,0}(z, \theta_0)}{\cos(\theta_R) - \sin(\theta_R)m_{-,0}(z, \theta_0)} \\
&= -\frac{\sin(\theta_R) + \cos(\theta_R)u'_{-, \theta_0}(z, R)}{\cos(\theta_R) - \sin(\theta_R)u'_{-, \theta_0}(z, R)} \\
&= \frac{1}{\sin^2(\theta_R)} [G_{\theta_0, \theta_R}(z, R, R) + \sin(\theta_R) \cos(\theta_R)], \tag{6.4}
\end{aligned}$$

$$\theta_0 \in [0, 2\pi), \theta_R \in [0, 2\pi) \setminus \{0, \pi\},$$

$$\begin{aligned}
\Lambda_{0,0}(z)_{2,2} &= -m_{-,0}(z, 0) \\
&= \lim_{0 < x < x', x \uparrow R} \partial_x \partial_{x'} G_{0,0}(z, x, x'). \tag{6.5}
\end{aligned}$$

Here we used (3.29) and (cf. also (3.28))

$$\begin{aligned}
G_{\theta_0, \theta_R}(z, R, R) &= \frac{-\sin(\theta_R)}{\cos(\theta_R) - \sin(\theta_R)m_{-,0}(z, \theta_0)} \\
&= \sin(\theta_R) [-\cos(\theta_R) - \sin(\theta_R)m_{-, \theta_R}(z, \theta_0)]. \tag{6.6}
\end{aligned}$$

Similarly, the off-diagonal terms of $\Lambda_{\theta_0, \theta_R}(z)$ in (3.21) can be written as

$$\begin{aligned} \Lambda_{\theta_0, \theta_R}(z)_{1,2} &= \Lambda_{\theta_0, \theta_R}(z)_{2,1} \\ &= \frac{1}{\sin(\theta_R)} G_{\theta_0, \theta_R}(z, R, R) \begin{cases} \frac{-u'_{-, \theta_0}(z, 0)}{\cos(\theta_0)}, & \theta_0 \in [0, 2\pi] \setminus \{\pi/2, 3\pi/2\}, \\ \frac{u_{-, \theta_0}(z, 0)}{\sin(\theta_0)}, & \theta_0 \in [0, 2\pi] \setminus \{0, \pi\}, \\ \theta_R \in [0, 2\pi] \setminus \{0, \pi\}, \end{cases} \end{aligned} \quad (6.7)$$

$$\begin{aligned} &= \frac{1}{\sin(\theta_0)} G_{\theta_0, \theta_R}(z, 0, 0) \begin{cases} \frac{u'_{+, \theta_R}(z, R)}{\cos(\theta_R)}, & \theta_R \in [0, 2\pi] \setminus \{\pi/2, 3\pi/2\}, \\ \frac{u_{+, \theta_R}(z, R)}{\sin(\theta_R)}, & \theta_R \in [0, 2\pi] \setminus \{0, \pi\}, \\ \theta_0 \in [0, 2\pi] \setminus \{0, \pi\}. \end{cases} \end{aligned} \quad (6.8)$$

Next we turn to (variants of) the 2×2 Weyl–Titchmarsh matrix corresponding to H_{θ_0, θ_R} with respect to an interior reference point $x_0 \in (0, R)$. We start by introducing (again, choosing $z \in \mathbb{C} \setminus \mathbb{R}$ for notational simplicity)

$$m_{+,0}(z, x_0, \theta_R) = \frac{u'_{+, \theta_R}(z, x_0)}{u_{+, \theta_R}(z, x_0)}, \quad x_0 \in (0, R), \quad (6.9)$$

$$m_{-,0}(z, x_0, \theta_0) = \frac{u'_{-, \theta_0}(z, x_0)}{u_{-, \theta_0}(z, x_0)}, \quad x_0 \in (0, R), \quad (6.10)$$

and more generally,

$$m_{+, \alpha}(z, x_0, \theta_R) = \frac{-\sin(\alpha) + \cos(\alpha)m_{+,0}(z, x_0, \theta_R)}{\cos(\alpha) + \sin(\alpha)m_{+,0}(z, x_0, \theta_R)}, \quad \alpha \in [0, \pi), \quad (6.11)$$

$$m_{-, \alpha}(z, x_0, \theta_0) = \frac{-\sin(\alpha) + \cos(\alpha)m_{-,0}(z, x_0, \theta_0)}{\cos(\alpha) + \sin(\alpha)m_{-,0}(z, x_0, \theta_0)}, \quad \alpha \in [0, \pi). \quad (6.12)$$

We note that $m_{+, \alpha}(\cdot, x_0, \theta_R)$ and $-m_{-, \alpha}(\cdot, x_0, \theta_0)$ are known to be Herglotz functions (cf., e.g., [23, Sect. 9.5]), [57, Sect. II.8], [68, Sect. 6.5], [80, Ch. III]).

Associated with (6.9)–(6.12) one then defines the 2×2 Weyl–Titchmarsh matrices (the fundamental ingredient for inverse spectral theory for operators of the type H_{θ_0, θ_R}) by,

$$\begin{aligned} M_0(z, x_0, \theta_0, \theta_R) &= [M_{0,j,k}(z, x_0, \theta_0, \theta_R)]_{j,k=1,2} \\ &= [m_{-,0}(z, x_0, \theta_0) - m_{+,0}(z, x_0, \theta_R)]^{-1} \\ &\quad \times \begin{bmatrix} 1 & \frac{1}{2}[m_{-,0}(z, x_0, \theta_0) + m_{+,0}(z, x_0, \theta_R)] \\ \frac{1}{2}[m_{-,0}(z, x_0, \theta_0) + m_{+,0}(z, x_0, \theta_R)] & m_{-,0}(z, x_0, \theta_0)m_{+,0}(z, x_0, \theta_R) \end{bmatrix}, \end{aligned} \quad (6.13)$$

and more generally, by

$$\begin{aligned}
M_\alpha(z, x_0, \theta_0, \theta_R) &= [M_{\alpha,j,k}(z, x_0, \theta_0, \theta_R)]_{j,k=1,2} \\
&= [m_{-, \alpha}(z, x_0, \theta_0) - m_{+, \alpha}(z, x_0, \theta_R)]^{-1} \\
&\quad \times \begin{bmatrix} 1 & \frac{1}{2}[m_{-, \alpha}(z, x_0, \theta_0) + m_{+, \alpha}(z, x_0, \theta_R)] \\ \frac{1}{2}[m_{-, \alpha}(z, x_0, \theta_0) + m_{+, \alpha}(z, x_0, \theta_R)] & m_{-, \alpha}(z, x_0, \theta_0)m_{+, \alpha}(z, x_0, \theta_R) \end{bmatrix}, \\
&\hspace{15em} \alpha \in [0, \pi).
\end{aligned} \tag{6.14}$$

By inspection,

$$\det(M_\alpha(z, x_0, \theta_0, \theta_R)) = -1/4, \quad \alpha \in [0, \pi), \quad z \in \mathbb{C} \setminus \mathbb{R}, \tag{6.15}$$

$$\operatorname{Im}(M_\alpha(z, x_0, \theta_0, \theta_R)) > 0, \quad \alpha \in [0, \pi), \quad z \in \mathbb{C}_+. \tag{6.16}$$

$M_\alpha(\cdot, x_0, \theta_0, \theta_R)$ is thus a 2×2 matrix-valued Herglotz function and $M_{\alpha,j,j}(\cdot, x_0, \theta_0, \theta_R)$ are Herglotz functions for $j = 1, 2$. In particular, the matrices $M_\alpha(z, x_0, \theta_0, \theta_R)$ can be shown to have Herglotz representations of the type (4.31).

To connect $M_0(z, x_0, \theta_0, \theta_R)$ and $M_\alpha(z, x_0, \theta_0, \theta_R)$ with the Green's function $G_{\theta_0, \theta_R}(z, x, x')$ of H_{θ_0, θ_R} , we first introduce a bit of notation:

$$\begin{aligned}
\partial_1 G_{\theta_0, \theta_R}(z, x_0, x') &= \partial_{x_1} G_{\theta_0, \theta_R}(z, x_1, x') \Big|_{x_1=x_0}, \\
\partial_2 G_{\theta_0, \theta_R}(z, x, x_0) &= \partial_{x_2} G_{\theta_0, \theta_R}(z, x, x_2) \Big|_{x_2=x_0}, \\
\partial_1 \partial_2 G_{\theta_0, \theta_R}(z, x_0, x_0) &= \partial_{x_1} \partial_{x_2} G_{\theta_0, \theta_R}(z, x_1, x_2) \Big|_{x_1=x_0, x_2=x_0}, \quad \text{etc.}
\end{aligned} \tag{6.17}$$

The expressions (6.13) and (6.14) for $M_0(z, x_0, \theta_0, \theta_R)$ and $M_\alpha(z, x_0, \theta_0, \theta_R)$ then can be rewritten as follows:

$$M_{0,1,1}(z, x_0, \theta_0, \theta_R) = G_{\theta_0, \theta_R}(z, x_0, x_0). \tag{6.18}$$

$$\begin{aligned}
M_{0,1,2}(z, x_0, \theta_0, \theta_R) &= M_{0,2,1}(z, x_0, \theta_0, \theta_R) \\
&= (1/2)(\partial_1 + \partial_2)G_{\theta_0, \theta_R}(z, x_0 \pm 0, x_0 \mp 0),
\end{aligned} \tag{6.19}$$

$$M_{0,2,2}(z, x_0, \theta_0, \theta_R) = \partial_1 \partial_2 G_{\theta_0, \theta_R}(z, x_0, x_0), \tag{6.20}$$

and

$$\begin{aligned}
M_{\alpha,1,1}(z, x_0, \theta_0, \theta_R) \\
&= (\cos(\alpha) + \sin(\alpha)\partial_1)(\cos(\alpha) + \sin(\alpha)\partial_2)G_{\theta_0, \theta_R}(z, x_0, x_0).
\end{aligned} \tag{6.21}$$

$$\begin{aligned}
M_{\alpha,1,2}(z, x_0, \theta_0, \theta_R) &= M_{\alpha,2,1}(z, x_0, \theta_0, \theta_R) \\
&= (1/2)((\cos(\alpha) + \sin(\alpha)\partial_1)(-\sin(\alpha) + \cos(\alpha)\partial_2) \\
&\quad + (-\sin(\alpha) + \cos(\alpha)\partial_1)(\cos(\alpha) + \sin(\alpha)\partial_2))G_{\theta_0, \theta_R}(z, x_0 \pm 0, x_0 \mp 0),
\end{aligned} \tag{6.22}$$

$$\begin{aligned}
M_{\alpha,2,2}(z, x_0, \theta_0, \theta_R) \\
&= (-\sin(\alpha) + \cos(\alpha)\partial_1)(-\sin(\alpha) + \cos(\alpha)\partial_2)G_{\theta_0, \theta_R}(z, x_0, x_0).
\end{aligned} \tag{6.23}$$

For relevant references in the context of (6.9)–(6.23), we refer, for instance, to [20, Ch. III], [23, Ch. 9], [31, App. J], [32], [40], [41], [56, Ch. 2], [57, Ch. 2], [68, Ch. 6], [80, Chs. II, III], and the references cited therein.

A comparison of equations (6.1), (6.2), (6.4), (6.5), (6.7), (6.8) with equations (6.18)–(6.20) and (6.21)–(6.23), respectively, clearly exhibits the different character of $M_\alpha(z, x_0, \theta_0, \theta_R)$, $\alpha \in [0, \pi)$, and $\Lambda_{\theta_0, \theta_R}(z)$ despite the fact that both are 2×2 matrix-valued Herglotz functions whose associated matrix-valued measures contain all spectral information on H_{θ_0, θ_R} . Additional differences are highlighted in Remark 3.3, and we feel that $\Lambda_{\theta_0, \theta_R}(z)$ (and more generally, $\Lambda_{\theta_0, \theta_R}^{\theta'_0, \theta'_R}(z)$) is deserving of a more detailed study.

We conclude with a final observation:

Remark 6.1. *With only minor modifications, all results in this paper extend to general, regular three-coefficient differential expressions of the type*

$$\frac{1}{r} \left(-\frac{d}{dx} p \frac{d}{dx} + q \right), \quad x \in [0, R], \quad (6.24)$$

where

$$p > 0, r > 0 \text{ a.e. on } (0, R), \quad \frac{1}{p}, q, r \in L^1((0, R); dx), \quad (6.25)$$

generating Sturm–Liouville operator realizations in $L^2((0, R); r dx)$. One just needs to consistently replace the derivative f' for elements in operator domains and Wronskians by the first quasi-derivative (pf') .

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References

- [1] N. I. Akhiezer, I. M. Glazman. *Theory of linear operators in Hilbert space*, Volume II. Pitman, Boston, 1981.
- [2] S. Albeverio, J. F. Brasche, M. M. Malamud, H. Neidhardt. *Inverse spectral theory for symmetric operators with several gaps: scalar-type Weyl functions*. *J. Funct. Anal.*, 228 (2005), 144–188.
- [3] D. Alpay, J. Behrndt. *Generalized Q -functions and Dirichlet-to-Neumann maps for elliptic differential operators*. *J. Funct. Anal.*, 257 (2009), 1666–1694.

- [4] W. O. Amrein, D. B. Pearson. *M operators: a generalisation of Weyl–Titchmarsh theory*. J. Comp. Appl. Math., 171 (2004), 1–26.
- [5] Yu. M. Arlinskiĭ, E. R. Tsekanovskiĭ. *On von Neumann’s problem in extension theory of nonnegative operators*. Proc. Amer. Math. Soc., 131 (2003), 3143–3154.
- [6] Yu. M. Arlinskiĭ, E. R. Tsekanovskiĭ. *The von Neumann problem for nonnegative symmetric operators*. Integr. Eq. Operator Th., 51 (2005), 319–356.
- [7] S. A. Avdonin, M. I. Belishev, S. A. Ivanov. *Boundary control and a matrix inverse problem for the equation $u_{tt} - u_{xx} + V(x)u = 0$* . Math. USSR Sbornik, 72 (1992), 287–310.
- [8] S. Avdonin, P. Kurasov. *Inverse problems for quantum trees*. Inverse Probl. Imaging, 2 (2008), 1–21.
- [9] S. Avdonin, S. Lenhart, V. Protopopescu. *Solving the dynamical inverse problem for the Schrödinger equation by the boundary control method*. Inverse Probl., 18 (2002), 349–361.
- [10] S. Avdonin, S. Lenhart, V. Protopopescu. *Determining the potential in the Schrödinger equation from the Dirichlet to Neumann map by the boundary control method*. J. Inv. Ill-Posed Probl., 13 (2005), 1–14.
- [11] J. Behrndt, M. Langer. *Boundary value problems for partial differential operators on bounded domains*. J. Funct. Anal., 243 (2007), 536–565.
- [12] J. Behrndt, M. M. Malamud, H. Neidhardt. *Scattering matrices and Weyl functions*. Proc. London Math. Soc., 97 (2008), No. 3, 568–598.
- [13] J. F. Brasche, M. M. Malamud, H. Neidhardt. *Weyl functions and singular continuous spectra of self-adjoint extensions* in Stochastic processes, physics and geometry: New interplays. II. A volume in honor of Sergio Albeverio. F. Gesztesy, H. Holden, J. Jost, S. Paycha, M. Röckner, S. Scarlatti (eds.). Canadian Mathematical Society Conference Proceedings, Vol. 29, Amer. Math. Soc., Providence, RI, 2000, pp. 75–84.
- [14] J. F. Brasche, M. M. Malamud, H. Neidhardt. *Weyl function and spectral properties of self-adjoint extensions*. Integral Eq. Operator Th., 43 (2002), 264–289.
- [15] B. M. Brown, G. Grubb, I. G. Wood. *M-functions for closed extensions of adjoint pairs of operators with applications to elliptic boundary problems*. Math. Nachr., 282 (2009), 314–347.
- [16] M. Brown, J. Hinchcliffe, M. Marletta, S. Naboko, I. Wood. *The abstract Titchmarsh–Weyl M-function for adjoint operator pairs and its relation to the spectrum*. Integral Equ. Operator Th., 63 (2009), 297–320.
- [17] B. M. Brown, M. Marletta. *Spectral inclusion and spectral exactness for PDE’s on exterior domains*. IMA J. Numer. Anal., 24 (2004), 21–43.

- [18] B. M. Brown, M. Marletta, S. Naboko, I. Wood. *Boundary triplets and M -functions for non-selfadjoint operators, with applications to elliptic PDEs and block operator matrices*. J. London Math. Soc., 77 (2008), No. 2, 700–718.
- [19] J. Brüning, V. Geyley, K. Pankrashkin. *Spectra of self-adjoint extensions and applications to solvable Schrödinger operators*. Rev. Math. Phys., 20 (2008), 1–70.
- [20] R. Carmona, J. Lacroix. *Spectral theory of random Schrödinger operators*. Birkhäuser, Basel, 1990.
- [21] S. Clark, F. Gesztesy. *Weyl–Titchmarsh M -function asymptotics and Borg-type theorems for Dirac operators*. Trans. Amer. Math. Soc., 354 (2002), 3475–3534.
- [22] S. Clark, F. Gesztesy. *On self-adjoint and J -self-adjoint Dirac-type operators: A case study*, in Recent advances in differential equations and mathematical physics. N. Chernov, Y. Karpeshina, I. W. Knowles, R. T. Lewis, R. Weikard (eds.). Contemp. Math., Vol. 412, Amer. Math. Soc., Providence, RI, 2006, pp. 103–140.
- [23] E. A. Coddington, N. Levinson. *Theory of ordinary differential equations*. Krieger, Malabar, 1985.
- [24] V. A. Derkach, S. Hassi, M. M. Malamud, H. S. V. de Snoo. *Generalized resolvents of symmetric operators and admissibility*. Meth. Funct. Anal. Top., 6 (2000), No. 3, 24–55.
- [25] V. Derkach, S. Hassi, M. Malamud, H. de Snoo. *Boundary relations and their Weyl families*. Trans. Amer. Math. Soc., 358 (2006), 5351–5400.
- [26] V. A. Derkach, M. M. Malamud. *Generalized resolvents and the boundary value problems for Hermitian operators with gaps*. J. Funct. Anal., 95 (1991), 1–95.
- [27] V. A. Derkach, M. M. Malamud. *Characteristic functions of almost solvable extensions of Hermitian operators*. Ukrain. Math. J., 44 (1992), 379–401.
- [28] V. A. Derkach, M. M. Malamud. *The extension theory of Hermitian operators and the moment problem*. J. Math. Sci., 73 (1995), 141–242.
- [29] N. Dunford, J. T. Schwartz. *Linear operators Part II: Spectral theory*. Interscience, New York, 1988.
- [30] C. Fox, V. Oleinik, B. Pavlov. *A Dirichlet-to-Neumann map approach to resonance gaps and bands of periodic networks*, in Recent advances in differential equations and mathematical physics. N. Chernov, Y. Karpeshina, I. W. Knowles, R. T. Lewis, R. Weikard (eds.). Contemp. Math. Vol. 412, Amer. Math. Soc., Providence, RI, 2006, pp. 151–169.
- [31] F. Gesztesy, H. Holden. *Soliton equations and their algebro-geometric solutions. Volume I: $(1 + 1)$ -Dimensional continuous models*. Cambridge Studies in Advanced Mathematics, Vol. 79, Cambridge University Press, Cambridge, 2003.

- [32] F. Gesztesy, H. Holden, B. Simon, Z. Zhao. *Higher order trace relations for Schrödinger operators*. Rev. Math. Phys., 7 (1995), 893–922.
- [33] F. Gesztesy, N. J. Kalton, K. A. Makarov, E. Tsekanovskii. *Some applications of operator-valued Herglotz functions*, in Operator theory, system theory and related topics. The Moshe Livšic anniversary volume. D. Alpay, V. Vinnikov (eds.). Operator Theory: Advances and Applications, Vol. 123, Birkhäuser, Basel, 2001, pp. 271–321.
- [34] F. Gesztesy, K. A. Makarov, E. Tsekanovskii. *An Addendum to Krein’s formula*. J. Math. Anal. Appl., 222 (1998), 594–606.
- [35] F. Gesztesy, M. Mitrea. *Generalized Robin boundary conditions, Robin-to-Dirichlet maps, and Krein-type resolvent formulas for Schrödinger operators on bounded Lipschitz domains*, in Perspectives in partial differential equations, harmonic analysis and applications: A volume in honor of Vladimir G. Maz’ya’s 70th birthday. D. Mitrea, M. Mitrea (eds.). Proceedings of Symposia in Pure Mathematics, Vol. 79, Amer. Math. Soc., Providence, RI, 2008, pp. 105–173.
- [36] F. Gesztesy, M. Mitrea. *Robin-to-Robin maps and Krein-type resolvent formulas for Schrödinger operators on bounded Lipschitz domains*, in Modern analysis and applications. The Mark Krein centenary conference, Vol. 2. V. Adamyan, Y. M. Berezansky, I. Gohberg, M. L. Gorbachuk, V. Gorbachuk, A. N. Kochubei, H. Langer, G. Popov (eds.). Operator Theory: Advances and Applications, Vol. 191, Birkhäuser, Basel, 2009, pp. 81–113.
- [37] F. Gesztesy, M. Mitrea. *Nonlocal Robin Laplacians and some remarks on a paper by Filonov on eigenvalue inequalities*. J. Diff. Eq., 247 (2009), 2871–2896.
- [38] F. Gesztesy, M. Mitrea. *Self-adjoint extensions of the Laplacian and Krein-type resolvent formulas in nonsmooth domains*. Preprint, 2009.
- [39] F. Gesztesy, M. Mitrea, M. Zinchenko. *Variations on a theme of Jost and Pais*. J. Funct. Anal., 253 (2007), 399–448.
- [40] F. Gesztesy, R. Ratnaseelan, G. Teschl. *The KdV hierarchy and associated trace formulas*, in Recent developments in operator theory and its applications. I. Gohberg, P. Lancaster, and P. N. Shivakumar (eds.). Operator Theory: Advances and Applications, Vol. 87, Birkhäuser, Basel, 1996, pp. 125–163.
- [41] F. Gesztesy, B. Simon. *Uniqueness theorems in inverse spectral theory for one-dimensional Schrödinger operators*. Trans. Amer. Math. Soc., 348 (1996), 349–373.
- [42] F. Gesztesy, E. Tsekanovskii. *On matrix-valued Herglotz functions*. Math. Nachr., 218 (2000), 61–138.
- [43] F. Gesztesy and M. Zinchenko, *Boundary Data Maps, Perturbation Determinants, and Krein-Type Resolvent Formulas for Schrödinger Operators on Compact Intervals*, preprint, 2010.

- [44] V. I. Gorbachuk, M. L. Gorbachuk. *Boundary value problems for operator differential equations*. Kluwer, Dordrecht, 1991.
- [45] G. Grubb. *Krein resolvent formulas for elliptic boundary problems in nonsmooth domains*. *Rend. Semin. Mat. Univ. Politec. Torino*, 66 (2008), 271–297.
- [46] G. Grubb. *Distributions and operators*. Graduate Texts in Mathematics, Vol. 252, Springer, New York, 2009.
- [47] T. Kato. *Perturbation theory for linear operators*. Corr. printing of the 2nd ed., Springer, Berlin, 1980.
- [48] A. Kiselev, B. Simon. *Rank one perturbations with infinitesimal coupling*. *J. Funct. Anal.*, 130 (1995), 345–356.
- [49] M. G. Krein, I. E. Ovcharenko. *Q-functions and sc-resolvents of nondensely defined hermitian contractions*. *Sib. Math. J.*, 18 (1977), 728–746.
- [50] M. G. Krein, I. E. Ovčarenko. *Inverse problems for Q-functions and resolvent matrices of positive hermitian operators*. *Sov. Math. Dokl.*, 19 (1978), 1131–1134.
- [51] M. G. Krein, S. N. Saakjan. *Some new results in the theory of resolvents of hermitian operators*. *Sov. Math. Dokl.*, 7 (1966), 1086–1089.
- [52] M. G. Krein, Ju. L. Smul’jan. *On linear-fractional transformations with operator coefficients*. *Amer. Math. Soc. Transl.*, 103 (1974), No. 2, 125–152.
- [53] P. Kurasov. *Triplet extensions I: Semibounded operators in the scale of Hilbert spaces*. *J. Analyse Math.*, 107 (2009), 251–286.
- [54] P. Kurasov, S. T. Kuroda. *Krein’s resolvent formula and perturbation theory*. *J. Operator Th.*, 51 (2004), 321–334.
- [55] H. Langer, B. Textorius. *On generalized resolvents and Q-functions of symmetric linear relations (subspaces) in Hilbert space*. *Pacific J. Math.*, 72 (1977), 135–165.
- [56] B. M. Levitan. *Inverse Sturm–Liouville problems*. VNU Science Press, Utrecht, 1987.
- [57] B. M. Levitan, I. S. Sargsjan. *Introduction to spectral theory*. Amer. Math. Soc., Providence, RI, 1975.
- [58] M. M. Malamud, V. I. Mogilevskii. *Krein type formula for canonical resolvents of dual pairs of linear relations*. *Methods Funct. Anal. Topolog.*, 8 (2002), No. 4, 72–100.
- [59] V. A. Marchenko. *Some questions in the theory of one-dimensional linear differential operators of the second order, I*. *Trudy Moskov. Mat. Obšč.*, 1 (1952), 327–420. (Russian.) English transl. in *Amer. Math. Soc. Transl., Ser. 2*, 101 (1973), 1–104.

- [60] V. A. Marchenko. *Sturm–Liouville operators and applications*. Birkhäuser, Basel, 1986.
- [61] M. Marletta. *Eigenvalue problems on exterior domains and Dirichlet to Neumann maps*. J. Comp. Appl. Math., 171 (2004), 367–391.
- [62] S. Nakamura. *A remark on the Dirichlet–Neumann decoupling and the integrated density of states*. J. Funct. Anal., 179 (2001), 136–152.
- [63] G. Nenciu. *Applications of the Kreĭn resolvent formula to the theory of self-adjoint extensions of positive symmetric operators*. J. Operator Th., 10 (1983), 209–218.
- [64] K. Pankrashkin. *Resolvents of self-adjoint extensions with mixed boundary conditions*. Rep. Math. Phys., 58 (2006), 207–221.
- [65] B. Pavlov. *The theory of extensions and explicitly-soluble models*. Russ. Math. Surv., 42:6 (1987), 127–168.
- [66] B. Pavlov. *S-matrix and Dirichlet-to-Neumann operators*. Ch. 6.1.6 in *Scattering: Scattering and inverse scattering in pure and applied science*, Vol. 2. R. Pike, P. Sabatier (eds.). Academic Press, San Diego, 2002, pp. 1678–1688.
- [67] B. Pavlov. *Krein formula with compensated singularities for the ND-mapping and the generalized Kirchhoff condition at the Neumann Schrödinger junction*. Russ. J. Math. Phys., 15 (2008), 364–388.
- [68] D. B. Pearson. *Quantum scattering and spectral theory*. Academic Press, London, 1988.
- [69] A. Posilicano. *A Krein-like formula for singular perturbations of self-adjoint operators and applications*. J. Funct. Anal., 183 (2001), 109–147.
- [70] A. Posilicano. *Self-adjoint extensions by additive perturbations*. Ann. Scuola Norm. Sup. Pisa Cl. Sci. (5), 2 (2003), No. 1, 1–20.
- [71] A. Posilicano. *Boundary triples and Weyl functions for singular perturbations of self-adjoint operators*. Meth. Funct. Anal. Topology, 10 (2004), No. 2, 57–63.
- [72] A. Posilicano. *Self-adjoint extensions of restrictions*. Operators and Matrices, 2 (2008), 483–506.
- [73] A. Posilicano, L. Raimondi. *Krein’s resolvent formula for self-adjoint extensions of symmetric second-order elliptic differential operators*. J. Phys. A: Math. Theor., 42 (2009), 015204 (11pp).
- [74] A. Rybkin. *On the boundary control approach to inverse spectral and scattering theory for Schrödinger operators*. Inverse Probl. Imaging, 3 (2009), 139–149.
- [75] V. Ryzhov. *A general boundary value problem and its Weyl function*. Opuscula Math., 27 (2007), 305–331.

- [76] V. Ryzhov. *Weyl–Titchmarsh function of an abstract boundary value problem, operator colligations, and linear systems with boundary control*. *Complex Anal. Operator Theory*, 3 (2009), 289–322.
- [77] V. Ryzhov. *Spectral boundary value problems and their linear operators*. Preprint, 2009.
- [78] Sh. N. Saakjan. *On the theory of the resolvents of a symmetric operator with infinite deficiency indices*. *Dokl. Akad. Nauk Arm. SSR*, 44 (1965), 193–198. (Russian.)
- [79] A. V. Straus. *Extensions and generalized resolvents of a non-densely defined symmetric operator*. *Math. USSR Izv.*, 4 (1970), 179–208.
- [80] E. C. Titchmarsh. *Eigenfunction expansions, Part I*. 2nd ed., Clarendon Press, Oxford, 1962.
- [81] E. C. Titchmarsh. *The theory of functions*. 2nd ed., Oxford University Press, Oxford, 1985.
- [82] E. R. Tsekanovskii, Yu. L. Shmul’yan. *The theory of bi-extensions of operators on rigged Hilbert spaces. Unbounded operator colligations and characteristic functions*. *Russ. Math. Surv.*, 32:5 (1977), 73–131.