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CONCENTRATION OF SOLAR ENERGY USING A CASSEGRAIN TYPE SOLAR FURNACE

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Abstract

A solar furnace consisting of a paraboloid of revolution that tracks the sun, and a hyperboloid of revolution reflector that has a focus in common with the paraboloid is analyzed using a three-dimensional ray trace to determine the image shape and size, and for the concentration to be obtained for various eccentricities e of the hyperboloid when used with an 84 foot diameter paraboloid (radar dish).

INTRODUCTION

A schematic view of the problem to be considered is shown in Figure 1. The paraboloid tracks the sun, and the sun's image is reflected to the hyperboloidal surface, and then re-reflected to the focal plane.



Fig. 1 Cassegrain Geometry

THREE DIMENSIONAL RAY TRACE

Vector a, ray from the sun

A ray from any position on the sun may be given as

$$\bar{A} = A_1 \bar{i} + A_2 \bar{j} + A_3 \bar{k}$$
(1)
and

$$A = \sqrt{A_1^2 + A_2^2 + A_3^2}$$
(2)

SO

$$\bar{a} = \frac{\bar{A}}{\bar{A}} = a_1 \bar{I} + a_2 \bar{J} + a_3 \bar{k}$$
(3)

Equation (3) is the normal form of a ray from the sun.

Paraboloid of revolution (concentrator)

The equation of a paraboloid of revolution, see Figure 1, is

$$y = a_0 + (\frac{x^2 + z^2}{4f_0})$$
 (4)

where

a = distance from the vertex of the paraboloid to the origin

f = focal length of the paraboloid

and in Figure 1, $r = x^2 + z^2$. We may also write

$$\phi_{\rm p} = y - \frac{(x^2 + z^2)}{4f_{\rm o}} = a_{\rm o}$$
 (5)

Normal to the paraboloid, b

$$\nabla \phi_{\mathbf{p}} = \frac{\partial \phi}{\partial \mathbf{x}} \mathbf{p} \mathbf{i} + \frac{\partial \phi}{\partial \mathbf{y}} \mathbf{p} \mathbf{j} + \frac{\partial \phi}{\partial \mathbf{z}} \mathbf{p} \mathbf{k} = -\frac{2\mathbf{x}\mathbf{i}}{4\mathbf{f}_{0}} + \mathbf{j} - \frac{2\mathbf{z}\mathbf{k}}{4\mathbf{f}_{0}}$$
$$= \mathbf{\bar{B}} = \mathbf{B}_{1}\mathbf{i} + \mathbf{B}_{2}\mathbf{j} + \mathbf{E}_{3}\mathbf{k} \qquad (6)$$
and

$$B = \sqrt{B_1^2 + B_2^2 + B_3^2}$$
(7)

SO

Б

$$= \frac{\bar{B}}{\bar{B}} = b_1 \bar{i} + b_2 \bar{j} + b_3 \bar{k}$$
(8)

Angle 0 between a, b

From Figure 2, we see that
$$\vec{a} \cdot \vec{b} = ab \cos\theta = \cos\theta$$

$$0 = \cos^{-1} \left[a_1 b_1 + a_2 b_2 + a_3 b_3 \right]$$
 (1)



Fig. 2 Vectors at Surface of Paraboloid

Vector
$$\bar{p}$$
, normal to \bar{a} , \bar{b}
 $\bar{P} = \bar{a} \times \bar{b} = \begin{vmatrix} \bar{i} & \bar{j} & \bar{k} \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix}$
 $= P_1 \bar{i} + P_2 \bar{j} + P_3 \bar{k}$ (11)

and

$$P = \sqrt{P_1^2 + P_2^2 + P_3^2}$$
(12)

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$$\bar{p} = \frac{\bar{p}}{\bar{p}} = p_1 \bar{i} + p_2 \bar{j} + p_3 \bar{k}$$
 (13)

Plane ϕ_{ab} containing \bar{a} , \bar{b}

$$\phi_{ab} = p_1 x + p_2 y + p_3 z = d_{ab}$$
(14)

where

 $d_{ab} = p_1 x_0 + p_2 y_0 + p_3 z_0$ (15)

where (x_0, y_0, z_0) are the coordinates of any point of interest on the paraboloid.

Reflected ray c

The vector equations for determining c are:

 $\mathbf{\bar{b}} \cdot \mathbf{\bar{c}} = \mathbf{b}\mathbf{c} \cos\theta = \cos\theta$ (16) $\mathbf{\bar{a}} \cdot \mathbf{\bar{c}} = \mathbf{a}\mathbf{c} \cos2\theta = \cos2\theta$ (17)

$$\bar{p} \cdot \bar{c} = 0$$

Thus

$$\bar{\mathbf{c}} = \mathbf{c}_1 \bar{\mathbf{i}} + \mathbf{c}_2 \bar{\mathbf{j}} + \mathbf{c}_3 \bar{\mathbf{k}}$$
(19)

where

$$c_1 = \frac{D_1}{|a|}, c_2 = \frac{D_2}{|a|} \text{ and } c_3 = \frac{D_3}{|a|}$$
 (20)

(18)

(9)

and
$$|a| = \begin{vmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{vmatrix}$$
 (21)

and D_r , r = 1,2,3 is determined by replacing the <u>rth</u> column of |a| with a column of k's. Also

$$a_{11} = b_1$$
 $a_{21} = a_1$ $a_{31} = p_1$ $k_1 = \cos\theta$
 $a_{12} = b_2$ $a_{22} = a_2$ $a_{32} = p_2$ $k_2 = \cos2\theta$
 $a_{13} = b_3$ $a_{23} = a_3$ $a_{33} = p_3$ $k_3 = 0$

$$\bar{\mathbf{R}} = \bar{\mathbf{c}} \times \bar{\mathbf{p}} = \begin{vmatrix} \mathbf{I} & \mathbf{j} & \mathbf{\bar{k}} \\ \mathbf{c}_1 & \mathbf{c}_2 & \mathbf{c}_3 \\ \mathbf{p}_1 & \mathbf{p}_2 & \mathbf{p}_3 \end{vmatrix}$$

$$R_1 \overline{i} + R_2 \overline{j} + R_3 \overline{k}$$
 (22)

and

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$$R = \sqrt{R_1^2 + R_2^2 + R_3^2}$$
(23)

$$\bar{\mathbf{r}} = \frac{\bar{\mathbf{R}}}{\bar{\mathbf{R}}} = \mathbf{r}_1 \bar{\mathbf{i}} + \mathbf{r}_2 \bar{\mathbf{j}} + \mathbf{r}_3 \bar{\mathbf{k}}$$
(24)

Plane
$$\phi_{cp}$$
 containing \bar{c}, \bar{p}
 $\phi_{cp} = r_1 x + r_2 y + r_3 z = d_{cp}$ (25)

where

$$d_{cp} = r_1 x_0 + r_2 y_0 + r_3 z_0$$
(26)

Intersection of reflected ray \bar{c} with the hyperboloid

We previously developed the equations of two planes which contain \bar{c} , so we may write Equations (14) and (25) alternately as

$$p_1 x + p_3 z = d_{ab} - p_2 y$$
 (27)

$$r_1 x + r_3 z = d_{CD} - r_2 y$$
 (28)

and so

$$x = \frac{D_1}{|b|}, \quad z = \frac{D_2}{|b|}$$
 (29)

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where

$$|\mathbf{b}| = \begin{vmatrix} \mathbf{b}_{11} & \mathbf{b}_{12} \\ \mathbf{b}_{21} & \mathbf{b}_{22} \end{vmatrix}, \quad \mathbf{D}_{1} = \begin{vmatrix} \mathbf{a}_{1} & \mathbf{b}_{12} \\ \mathbf{a}_{2} & \mathbf{b}_{22} \end{vmatrix},$$

$$\mathbf{D}_{2} = \begin{vmatrix} \mathbf{b}_{11} & \mathbf{a}_{1} \\ \mathbf{b}_{21} & \mathbf{a}_{2} \end{vmatrix}$$
(30)

and where

 $b_{11} = p_1$ $b_{21} = r_1$ $\ell_1 = d_{ab} - p_2 y$ $b_{12} = p_3$ $b_{22} = r_3$ $\ell_2 = d_{cp} - r_2 y$

Thus we can write

$$\mathbf{x} = \mathbf{B}\mathbf{B}_1 + \mathbf{B}\mathbf{B}_2\mathbf{y} \tag{31}$$

and

 $z = BB_3 + BB_4 y \tag{32}$

where y is known, or assumed.

Hyperboloid of revolution (cassegrain reflector)

The equation for a hyperboloid of revolution, see Figure 1, is

$$\frac{y^2}{a^2} - (\frac{x^2 + z^2}{b^2}) = 1$$
(33)

where

$$F = ae = \frac{1}{2}$$
 distance between foci

and

$$e = \sqrt{\frac{a^2 + b^2}{a}} > 1 = eccentricity$$

So, using Equations (31) and (32), we may write alternately

$$F(y) = \frac{y^2}{a^2} - 1 - \frac{1}{b^2} [(BB_1 + BB_2 y)^2 + (BB_3 + BB_4 y)^2]$$
(34)

and

$$F'(y) = \frac{2y}{a^2} - \frac{2}{b^2} [(BB_1 + BB_2 y)BB_2 + (BB_3 + BB_4 y)BB_4]$$
(35)

Newton's method

A recursion method for finding y accurately, when an approximate value of y is known, is

$$\mathbf{y}_{n+1} = \mathbf{y}_n - \frac{\mathbf{F}(\mathbf{y}_n)}{\frac{\mathbf{F}(\mathbf{y}_n)}{\mathbf{F}(\mathbf{y}_n)}}$$
(36)

Thus starting with an estimate of y, (y_n) , a new value of y, (y_{n+1}) , is computed. The process is continued until the difference between y_{n+1} and y_n is as small as desired. Using this method, the coordinates of the pierce point of \bar{c} with the hyperboloid are found, namely (x_1, y_1, z_1) .

Normal to the hyperboloid, d

$$\phi_{n} = \frac{\partial \phi}{\partial \mathbf{x}^{n}} \mathbf{i}^{T} + \frac{\partial \phi}{\partial \mathbf{y}^{n}} \mathbf{j}^{T} + \frac{\partial \phi}{\partial \mathbf{z}^{n}} \mathbf{k}$$

$$= -\frac{2\mathbf{x}\mathbf{i}}{b^{2}} + \frac{2\mathbf{y}\mathbf{j}}{a^{2}} - \frac{2\mathbf{x}\mathbf{k}}{b^{2}} = \mathbf{D}$$

$$= D_{1}\mathbf{i}^{T} + D_{2}\mathbf{j}^{T} + D_{3}\mathbf{k} \qquad (37)$$

and

so

Π

$$D = \sqrt{D_1^2 + D_2^2 + D_3^2}$$
(38)

$$\bar{\mathbf{d}} = \frac{\bar{\mathbf{D}}}{\bar{\mathbf{D}}} = \mathbf{d}_1 \bar{\mathbf{i}} + \mathbf{d}_2 \bar{\mathbf{j}} + \mathbf{d}_3 \bar{\mathbf{k}}$$
(39)

Angle
$$\gamma$$
 between \overline{c} , \overline{d}
From Figure 3, we see that
 $\overline{c} \cdot \overline{d} = cd \cos \gamma = \cos \gamma$ (40)
so

$$\gamma = \cos^{-1} [c_1 d_1 + c_2 d_2 + c_3 d_3]$$
(41)



Fig. 3 Vectors at Surface of Hyperboloid

Vector t, normal to c, d

$$\bar{\mathbf{T}} = \bar{\mathbf{c}} \times \bar{\mathbf{d}} = \begin{vmatrix} \bar{\mathbf{i}} & \bar{\mathbf{j}} & \bar{\mathbf{k}} \\ \mathbf{c}_1 & \mathbf{c}_2 & \mathbf{c}_3 \\ \mathbf{d}_1 & \mathbf{d}_2 & \mathbf{d}_3 \end{vmatrix}$$

$$= \mathbf{T}_{1} \mathbf{\overline{i}} + \mathbf{T}_{2} \mathbf{\overline{j}} + \mathbf{T}_{3} \mathbf{\overline{k}}$$
(42)

and

$$T = \sqrt{T_1^2 + T_2^2 + T_3^2}$$
(43)

so

$$t = \frac{\overline{T}}{\overline{T}} = t_1 \overline{i} + t_2 \overline{j} + t_3 \overline{k}$$
(44)

Re-reflected ray e

The vector equations for determining $\bar{\mathbf{e}}$ are:

$$\bar{d} \cdot \bar{e} = de \cos \gamma = \cos \gamma$$
 (45)

$$c \cdot e = ce \cos 2\gamma = \cos 2\gamma$$
 (46)

$$t \cdot e = 0$$
 (47)

Thus

$$\bar{e} = e_1 \bar{i} + e_2 \bar{j} + e_3 \bar{k}$$
 (48)

where

$$e_1 = \frac{D_1}{|c|}, e_2 = \frac{D_2}{|c|} \text{ and } e_3 = \frac{D_3}{|c|}$$
 (49)

and

$$|c| = \begin{vmatrix} c_{11} & c_{12} & c_{13} \\ c_{21} & c_{22} & c_{23} \\ c_{31} & c_{32} & c_{33} \end{vmatrix}$$

and $D_r, r = 1, 2, 3$ is determined, as before by replacing the rth column of |c| by a column k's. Also

$$c_{11} = d_1 \qquad c_{21} = c_1 \qquad c_{31} = t_1 \qquad k_1 = \cos\gamma$$

$$c_{12} = d_2 \qquad c_{22} = c_2 \qquad c_{32} = t_2 \qquad k_2 = \cos 2\gamma$$

$$c_{13} = d_3 \qquad c_{23} = c_3 \qquad c_{33} = t_3 \qquad k_3 = 0$$

Vector s, normal to e, t

$$\overline{\mathbf{S}} = \overline{\mathbf{e}} \times \overline{\mathbf{t}} = \begin{vmatrix} \overline{\mathbf{i}} & \overline{\mathbf{j}} & \overline{\mathbf{k}} \\ \mathbf{e}_1 & \mathbf{e}_2 & \mathbf{e}_3 \\ \mathbf{t}_1 & \mathbf{t}_2 & \mathbf{t}_3 \end{vmatrix}$$

$$= s_1 \bar{i} + s_2 \bar{j} + s_3 \bar{k}$$
 (51)

and

$$S = \sqrt{S_1^2 + S_2^2 + S_3^2}$$
(52)

$$\overline{s} = \frac{\overline{s}}{\overline{s}} = s_1 \overline{1} + s_2 \overline{j} + s_3 \overline{k}$$
(53)

Plane d_{cd} containing c,d

$$\phi_{cd} = t_1 x + t_2 y + t_3 z = d_{cd}$$
 (54)

where

so

$$d_{cd} = t_1 x_1 + t_2 y_1 + t_3 z_1$$
(55)

and where (x_1, y_1, z_1) are the coordinates of the point of intersection of the ray \bar{c} with the hyperboloid.

Plane ϕ_{te} containing t,e

$$\phi_{te} = s_1 x + s_2 y + s_3 z = d_{te}$$
 (56)

where

$$d_{te} = s_1 x_1 + s_2 y_1 + s_3 z_1$$
 (57)

Plane ϕ_{F} containing the focal plane

$$\phi_{F} = f_{1}x + f_{2}y + f_{3}z$$

= $f_{2}y = y = d_{f}$ (58)

where

$$d_{f} = y_{2}$$
(59)

and y_2 is the coordinate of the focal plane.

Intersection of the re-reflected ray e with the focal plane

The controlling equations for finding the pierce of \bar{e} in the focal plane are:

$$s_1 x + s_2 y + s_3 z = d_{te}$$
 (60)

$$t_1 x + t_2 y + t_3 z = d_{cd}$$
 (61)

$$y = y_2 \tag{62}$$

SO

$$x = x_2 = \frac{D_1}{|d|}, \quad z = z_2 = \frac{D_2}{|d|}$$
 (63)

where

$$|\mathbf{d}| = \begin{vmatrix} \mathbf{d}_{11} & \mathbf{d}_{12} \\ \mathbf{d}_{21} & \mathbf{d}_{22} \end{vmatrix}, \quad \mathbf{D}_1 = \begin{vmatrix} \mathbf{a}_1 & \mathbf{d}_{12} \\ \mathbf{a}_2 & \mathbf{d}_{22} \end{vmatrix},$$

$$D_{2} = \begin{vmatrix} d_{11} & \ell_{1} \\ d_{21} & \ell_{2} \end{vmatrix}$$
(64)

and where

$$d_{11} = s_1$$
 $d_{21} = t_1$ $\ell_1 = d_{te} - s_{2y}$
 $d_{12} = s_3$ $d_{22} = t_3$ $\ell_2 = d_{cd} - t_{2y}$

Thus the pierce point (x_2, y_2, z_2) of the ray \overline{e} with the focal plane is now known.

Radius of the pierce point in the focal plane, r

$$r_{t} = \sqrt{r_{2}^{2} + z_{2}^{2}}$$
 (65)

Figures 4-8 show the image size and shape having been reflected from points on the paraboloid having radii from the axis of symmetry $r_p = 2,12,22,32$ and 42 feet, respectively. It is seen that the image shapes are not circular, but are not severely distorted from a circular shape. This is in contrast to the images cast by the sun, from points on the paraboloid, to a flat plate oriented perdendicular to the axis of symmetry and placed at the focus of the paraboloid at low angles of incidence.

Concentration ratio

Writing an identity for the energy arriving at the image, gives

$$\pi (r_{p}^{2} - r_{h}^{2}) S_{c}^{\rho} p^{\rho} h = \pi r_{tm}^{2} S_{t}$$
(66)

and defining

$$C_{av} = \frac{s_t}{s_c} = \frac{(r_p^2 - r_h^2)}{r_{tm}^2} \rho_p \rho_h$$
 (67)

where

- S = solar constant
- St = average flux across the image at the focal plane
- $r_{D} = radius$ of parabolal considered

r_h = radius of necessary hyperboloid

- p p h = hemispherical reflectivity of
 paraboloidal surface and hyper boloidal surface, respectively
 (averaged)
 - rt = radius of ray pierce point in the focal plane measured from the focal point of the hyperboloid
- r_{tm} = maximum radius of ray pierce point, considering the entire paraboloid

We may also write

$$C_{av} = C_{i} n_{e}$$
 (68)

where

$$C_{i} = (\frac{r_{p}^{2} - r_{h}^{2}}{r_{tm}^{2}}) = \text{ideal concentration}$$
(geometric) (69)

and

$$n_e = \rho_p \rho_h = combined reflectivityfactor (70)$$

The concentration above may be increased using a compound paraboloidal concentrator (or a cone can be used with less effectiveness)(2) according to

 $C_{cpc} = \frac{1}{\sin^2 \zeta} = \frac{1}{\operatorname{concentration by}}$ $C_{cpc} = \frac{1}{\sin^2 \zeta} = \frac{1}{\operatorname{concentration by}}$ $C_{cpc} = \frac{1}{\sin^2 \zeta} = \frac{1}{\operatorname{concentrator}}$ $C_{cpc} = \frac{1}{\sin^2 \zeta}$ $C_{cpc} = \frac{1}{\sin^2 \zeta}$ $C_{cpc} = \frac{1}{\sin^2 \zeta}$ $C_{cpc} = \frac{1}{\cos^2 \zeta}$ $C_{cpc} = \frac{1}{\sin^2 \zeta}$ $C_{cpc} = \frac{1}{\cos^2 \zeta}$ $C_{cpc} = \frac{1}{\sin^2 \zeta}$ $C_{cpc} = \frac{1}{\cos^2 \zeta}$

and

$$C_{aug} = C_{cpc}C_{av} = concentration augmented (72)$$

where from Figure 9

$$\zeta = \tan^{-1} \left(\frac{r_{t} + r_{h}}{F + y_{l}} \right)$$
(73)

$$C_{aug} = C_{cpc}C_{i}^{n}e =$$

$$\frac{(r_{p}^{2} - r_{h}^{2})}{r_{tm}^{2}} \frac{\rho_{p}\rho_{h}}{\sin^{2}\varsigma}$$
(74)

and

$$C_{aug_i} = C_{cpc}C_i = \frac{C_{aug}}{n_e} =$$

ideal augmented concentration (75)

Figure 10 shows a plot of the ideal concentration, C_i vs. the eccentricity e, for various values of CL, the distance from the vertix of the paraboloid to the focal plane (see Figure 1).

Figure 11 shows a plot of the ideal augmented concentration, C_{aug} vs. the eccentricity e, for various values of CL.

Figure 12 shows a plot of the necessary radius of the hyperboloid, r_h vs. the eccentricity e, for various values of CL.



Fig. 4 Image Shape, $r_p = 2$ feet



Fig. 5 Image Shape, $r_p = 12$ feet



Fig. 6 Image Shape, $r_p = 22$ feet







Fig. 8 Image Shape, $r_p = 42$ feet



Fig. 9 Hyperboloid and Image Geometry



Fig. 10 Ideal Concentration vs. Eccentricity



Fig. 11 Ideal Augmented Concentration vs. Eccentricity



Fig. 12 Hyperboloid Radius vs. Eccentricity

NOMENCLATURE

Ā, B, D, P, R Ī,Ī = vectors ā, b, c, d, ē, p,r,t,s = normalized vectors |a|, |b|, |c|, |d| = determinants a = paraboloid constant a,b = hyperboloid constants BB1, BB2 BB_3 , BB_4 = constants C_{av}, C_i, C_{aug}, CL = distance from the vertex of the paraboloid to the focal plane $d_{ab}, d_{cd}, d_{cd}, d_{cp}, d_{te}$ = constants D_r = determinant f_o = focal length of paraboloid e = eccentricity F = ae F(y), F(y) = functions of y i, j, k = unit vectors k_r, k_r = constants r = radius

- r_{h} = radius of hyperboloid
- r_p = radius of paraboloid
- r = radius of ray pierce point in the focal plane measured from the focal point of the hyperboloid
- rtm = maximum radius of any ray pierce point considering the entire paraboloid
- S_c = solar constant
- St = average flux across the image at
 the focal plane
- x,y,z = coordinates
 - α_s = angle subtended by the sun = 0.009322 radians
- $\gamma, \theta, \zeta = angles$
 - ∇ = gradient operator

 ${}^{\phi}ab'{}^{\phi}cd'{}^{\phi}cp'$ ${}^{\phi}te'{}^{\phi}F$ = planes

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BIOGRAPHIES

Dr. Cobble is Professor of Mechanical Engineering at New Mexico State University. He received his Ph.D. from the University of Michigan in 1958. He has published in the areas of heat transfer, fluid mechanics and solar energy.

Dr. Hull received his Ph.D. from the University of Toledo in 1975. He is an Associate Professor of Mechanical Engineering at New Mexico State University. He has published in the areas of fatigue and failure analysis, and solar furnace design.